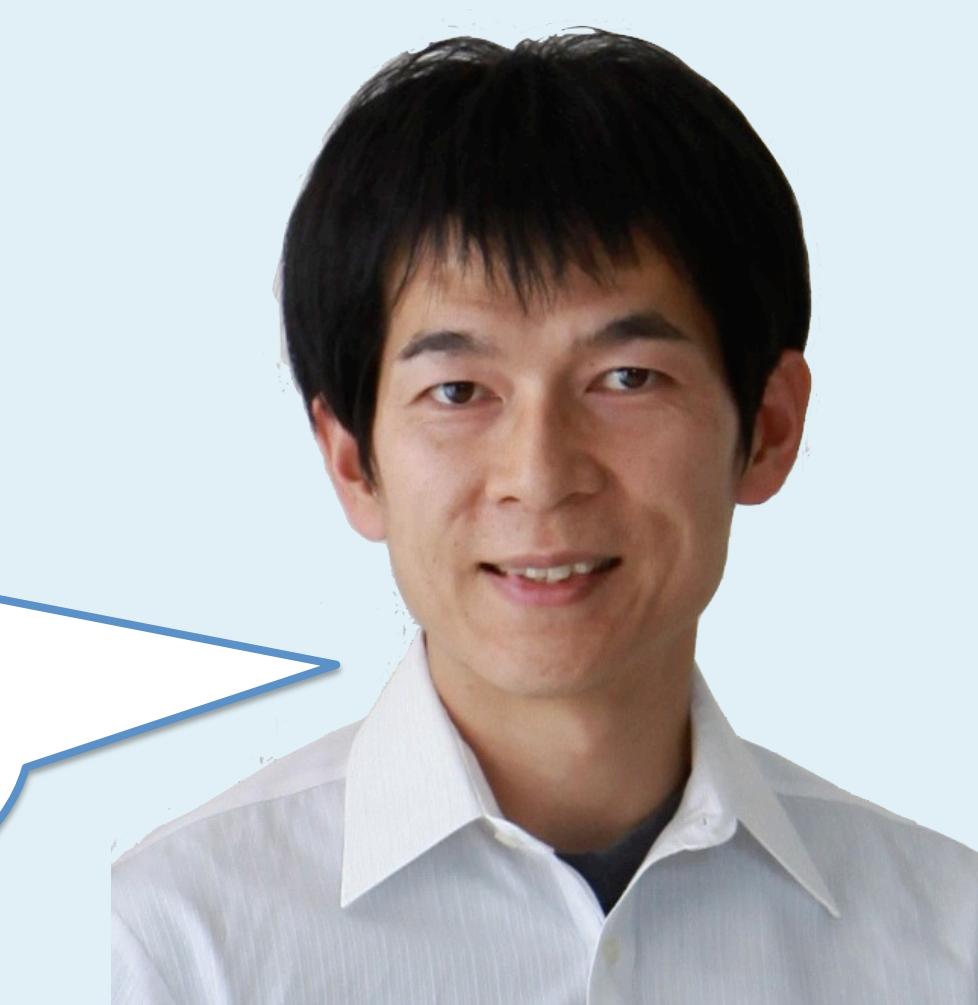


# Why the bolus velocity deserved to survive and how we use it?

## A 3D EP theory for all waves at all latitudes as given by (what we call) the impulse-bolus (IB) pseudomomentum

<http://www.jamstec.go.jp/frcgc/research/d1/aiki/>



- The IB pseudomomentum equation has previously been derived using a shallow-water model
- Pressure flux does not point in the direction of the group velocity of Rossby waves
- The IB flux points in the direction of the group velocity of all waves at all latitudes which is attributed to a divergence-form wave-induced pressure associated with the virial theorem

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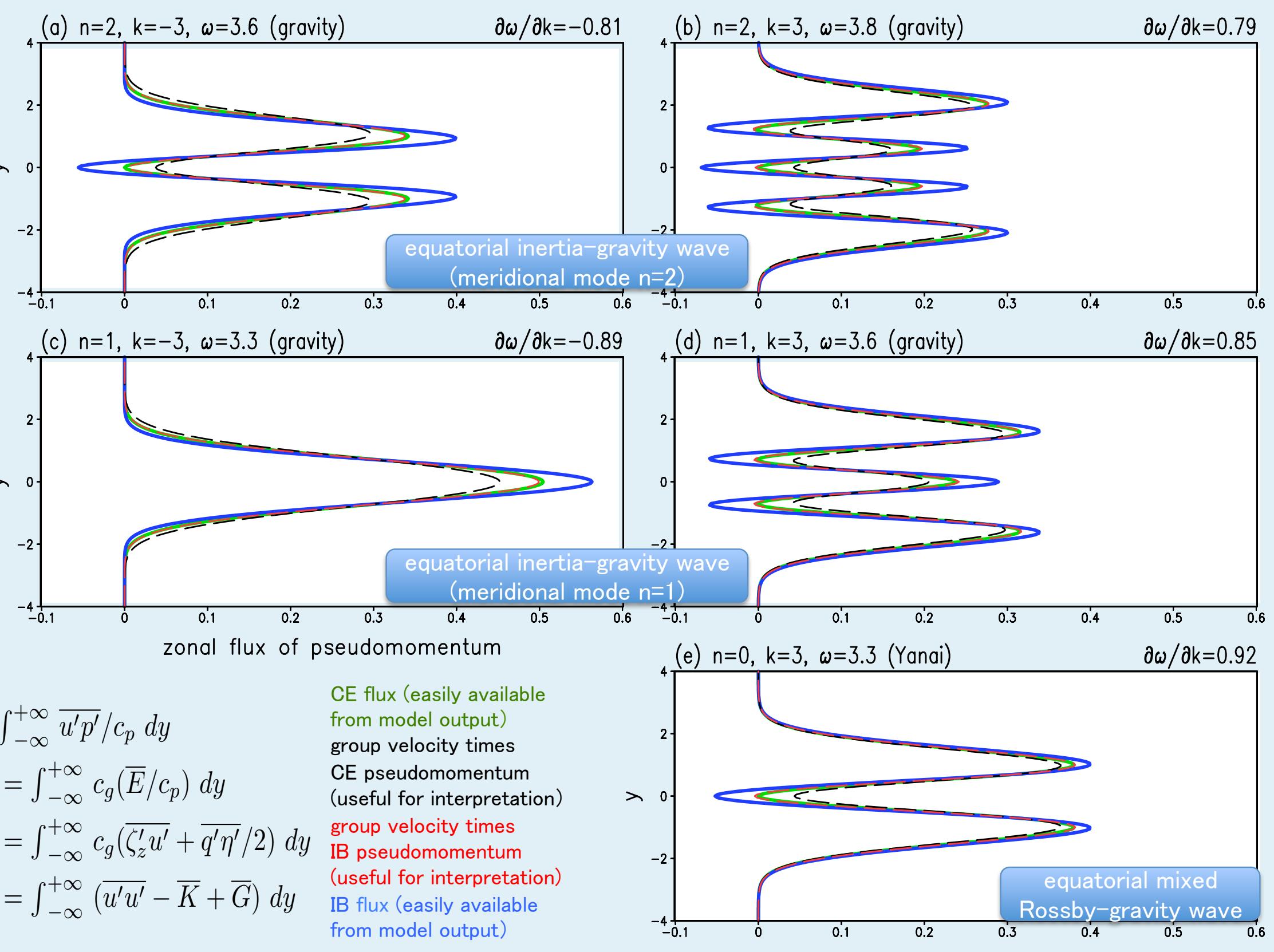
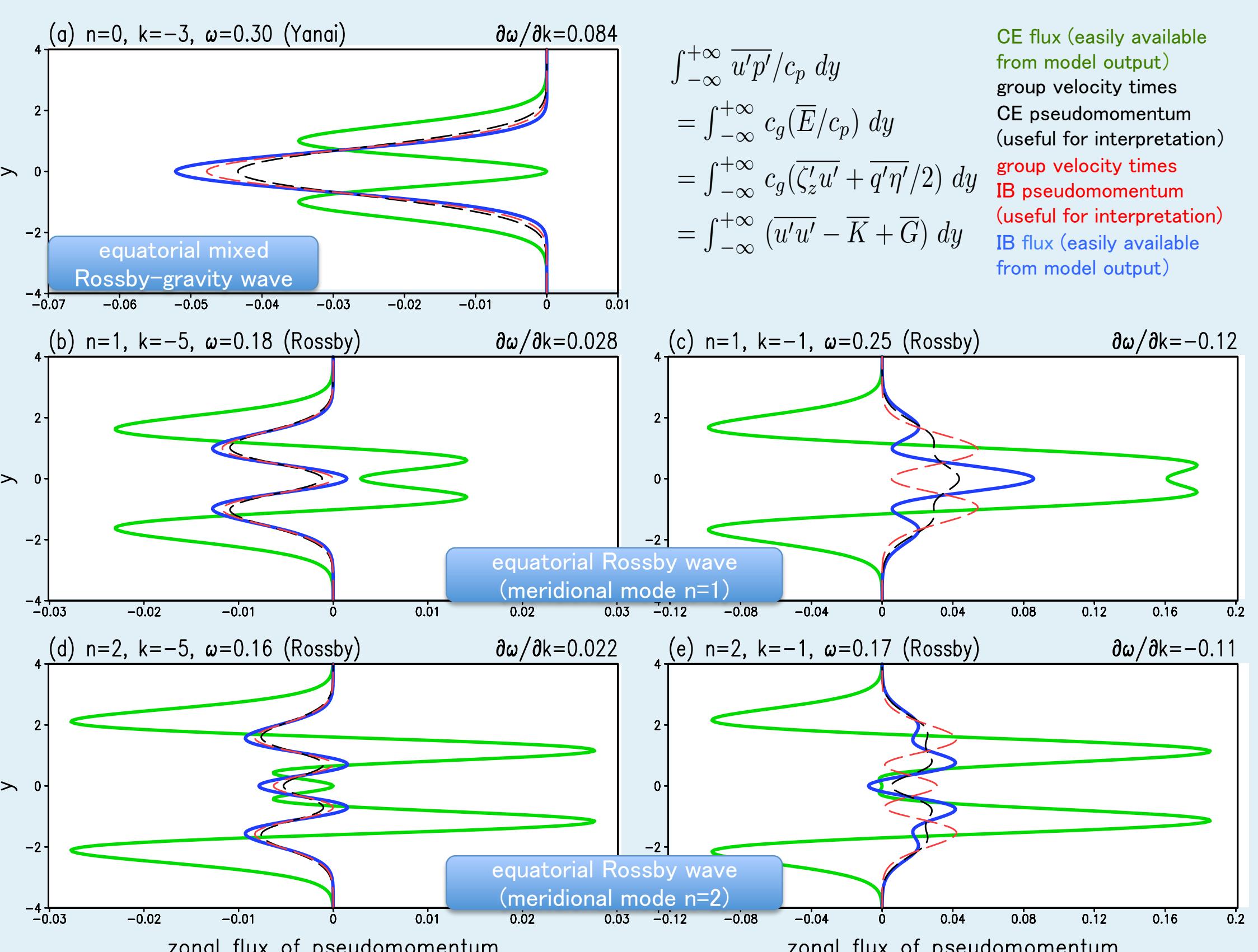
$$\begin{aligned} Q^{\dagger} &= \psi_{xx} + \psi_{yy} + (\psi_z f_0^2/N^2)_z \\ &(\partial_t + \bar{u}\partial_x)[\psi_{xx} + \psi_{yy} + (\psi_z f_0^2/N^2)_z] + \psi_x(\beta - \bar{u}_{yy}) = 0 \\ &\nabla_y z = \langle \partial_y, \partial_z \rangle \end{aligned}$$

deviation from Eulerian zonal mean  
 $A^{\dagger} = A - \bar{A}$   
 $\nabla_y z = \langle \partial_y, \partial_z \rangle$

QG Taylor-Bretherton identity (meridional flux of QGPV)  
 $\bar{Q}' \psi_x = -\nabla_y z \cdot \langle \bar{u}_x \bar{\psi}_y, -\bar{\psi}_x \bar{\psi}_z f_0^2/N^2 \rangle$   
 $\bar{v}'' u \bar{u} = -f_0 v'' \bar{u}^2/\bar{p}_z$

QG Eliassen-Palm relation (prognostic equation for the QG pseudomomentum)  
 $\partial_t [(\frac{1}{2}) \bar{Q}'^2 / (\bar{u}_{yy} - \beta)] + \nabla_y z \cdot \langle \bar{u}_x \bar{\psi}_y, -\bar{\psi}_x \bar{\psi}_z f_0^2/N^2 \rangle = 0$   
 $E/C_p^z$

Allows for the group velocity vector to be diagnosed from model output without performing a Fourier analysis



Governing equation for neutral waves at all latitudes:  
MIGWs (mid-latitude inertia-gravity waves)  
MRWs (mid-latitude Rossby waves)  
EQWs (all types of equatorial waves)

$$\begin{aligned} u'_t - (f_0 + \beta y)v' &= -p'_x, \\ v'_t + (f_0 + \beta y)u' &= -p'_y, \\ \rho'_t + w'\bar{p}_z &= 0, \\ u'_x + v'_y + w'_z &= 0, \\ \langle\langle u', v', w' \rangle\rangle &= \langle\langle \xi'_t, \eta'_t, \zeta'_t \rangle\rangle, \end{aligned}$$

$\zeta' \equiv -\rho'/\bar{p}_z = -p'_z/N^2$ ,  
 $K = (u'^2 + v'^2)/2$ ,  $G = (N^2/2)\zeta'^2$ ,  
 $q' \equiv v'_x - u'_y - (f_0 + \beta y)\zeta'_z$ ,  
 $q'_t + \beta v' = 0$ ,  
 $\eta' = -q'/\beta$ ,

Erte's potential vorticity

$$\begin{aligned} A' &= A - \bar{A} \\ \nabla &= \langle \partial_x, \partial_y, \partial_z \rangle \end{aligned}$$

in height coordinates and with overbar and single prime indicating Eulerian time mean and deviation from it

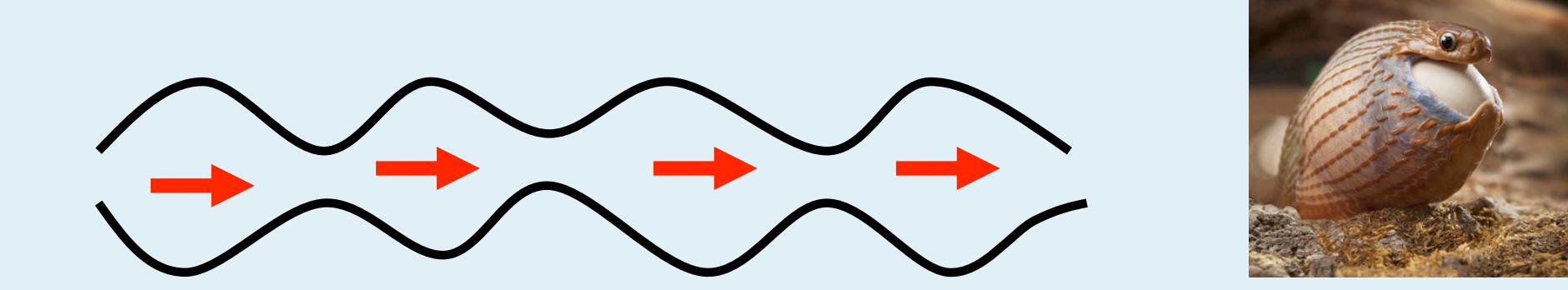
Approximated expressions using a Taylor expansion in height-coordinates

(overbar and single prime indicating an Eulerian time mean and deviation from it)

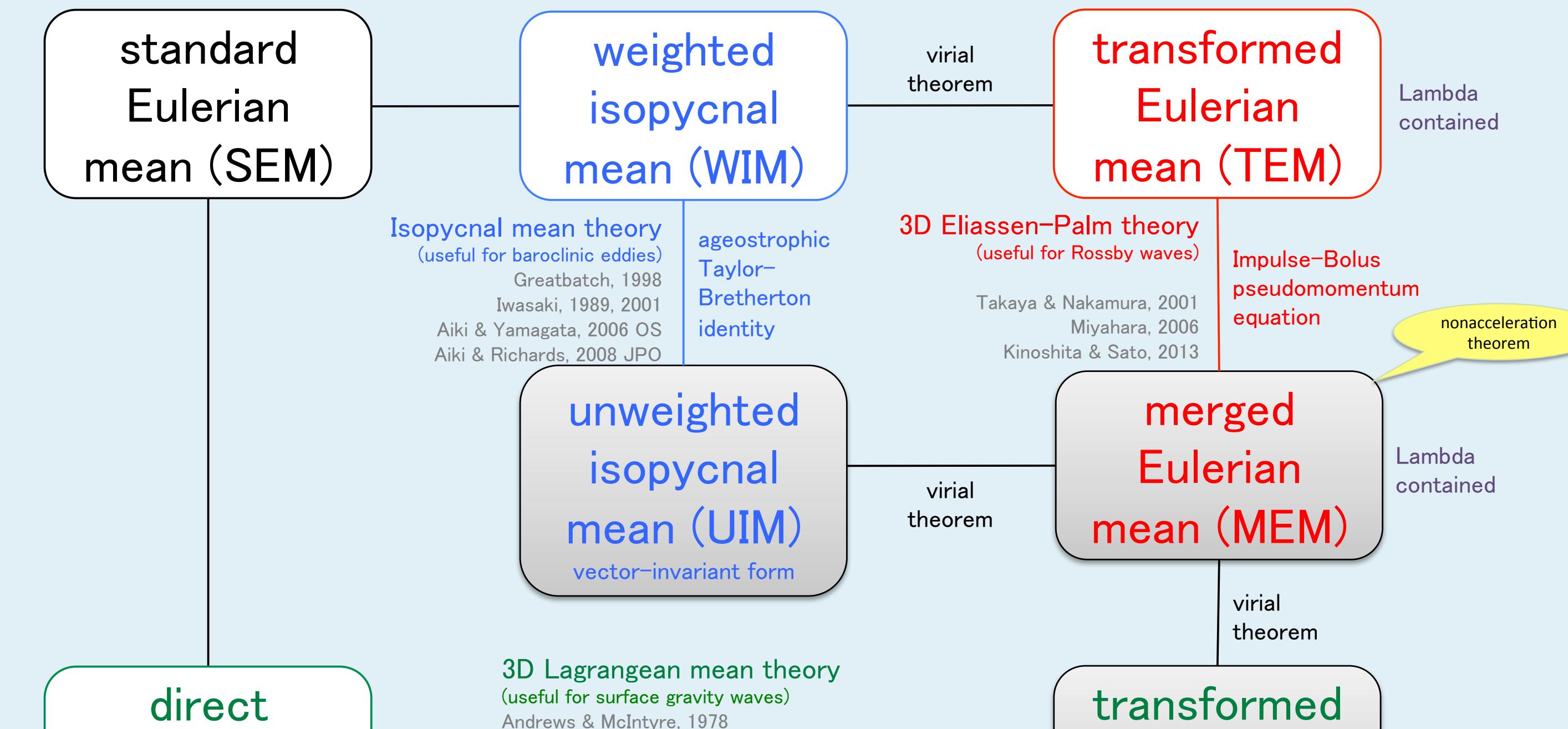
$$\begin{aligned} \bar{A}^L &= \bar{A} + \bar{\xi}' A'_x + \eta' A'_y + \zeta' A'_z && \text{3D Lagrangian mean} \\ \tilde{A} &= \bar{A} + \bar{\zeta}' A'_z && \text{Unweighted isopycnal mean} \\ \hat{A} &= \tilde{A} + \bar{\zeta}' A' = \bar{A} + (\bar{\zeta}' A')_z && \text{Thickness-weighted isopycnal mean} \end{aligned}$$

$u^{Stokes} \equiv \bar{u}^L - \bar{u} = \bar{\xi}' u'_x + \eta' u'_y + \zeta' u'_z$   
 $u^{qs} \equiv \hat{u} - \bar{u} = (\zeta' u')_z$   
 $u^{bolus} \equiv \hat{u} - \tilde{u} = \bar{\zeta}' u'$

Stokes-drift velocity  
quasi-Stokes velocity  
bolus velocity (Rhines, 1982)



Relationships between the low-pass filtered momentum equations in different theories for representing the effect of waves on the mean flow



$$\begin{aligned} (K+G)_t + \nabla \cdot \langle\langle u' p', v' p', w' p' \rangle\rangle &= 0, \quad c_p \equiv \omega/k, \quad c_g \equiv \partial\omega/\partial k \\ (E/c_p)_t + \nabla \cdot \langle\langle u' p'/c_p, v' p'/c_p, w' p'/c_p \rangle\rangle &= 0, \quad \bar{u}' p'/c_p \neq c_g (\bar{E}/c_p) \end{aligned}$$

CE flux (easily available from model output)  
group velocity times CE pseudomomentum (useful for interpretation)  
group velocity times IB pseudomomentum (useful for interpretation)  
IB flux (easily available from model output)

Prescription for the phase velocity in the denominator to be rewritten

See green and black lines in the left figure

$$\begin{aligned} \pi' &\equiv \int^t p' dt, \\ u' - f\eta' &= -\pi'_x, \\ v' + f\xi' &= -\pi'_y, \\ E &\equiv K + G \\ &= (u'^2 + v'^2 + N^2\zeta'^2)/2 \\ &= (u' \xi'_t + v' \eta'_t - \zeta' \pi'_{zt})/2, \end{aligned}$$

CE pseudomomentum

$$\begin{aligned} &[(-u' \xi'_x - v' \eta'_x + \zeta' \pi'_{zx})/2]_t \\ &= -\nabla \cdot \langle\langle -(\xi'_x p' + u' \pi'_x)/2, -(\eta'_x p' + v' \pi'_x)/2, -(\zeta'_x p' + w' \pi'_x)/2 \rangle\rangle, \end{aligned}$$

Redirecting the CE flux by singling out

$$\Lambda \equiv [(\xi' p'_x)_x + (\eta' p'_y)_y + (\zeta' p'_z)_z]/2$$

CE pseudomomentum

$$\begin{aligned} &[(-u' \xi'_x - v' \eta'_x + \zeta' \pi'_{zx})/2]_t \\ &= -\nabla \cdot \langle\langle (\xi' p'_x - u' \pi'_x)/2 - \Lambda, (\eta' p'_x - v' \pi'_x)/2, (\zeta' p'_x - w' \pi'_x)/2 \rangle\rangle, \end{aligned}$$

CE pseudomomentum

$$\begin{aligned} &[(-u' \xi'_x - v' \eta'_x + \zeta' \pi'_{zx})/2]_t \\ &= -\nabla \cdot \langle\langle E - v' v' + (v' \eta')_t/2, v' u' - (u' \eta')_t/2, \zeta' p'_x - (\zeta' \pi'_x)_t/2 \rangle\rangle, \end{aligned}$$

Impulse-Bolus (IB) pseudomomentum equation

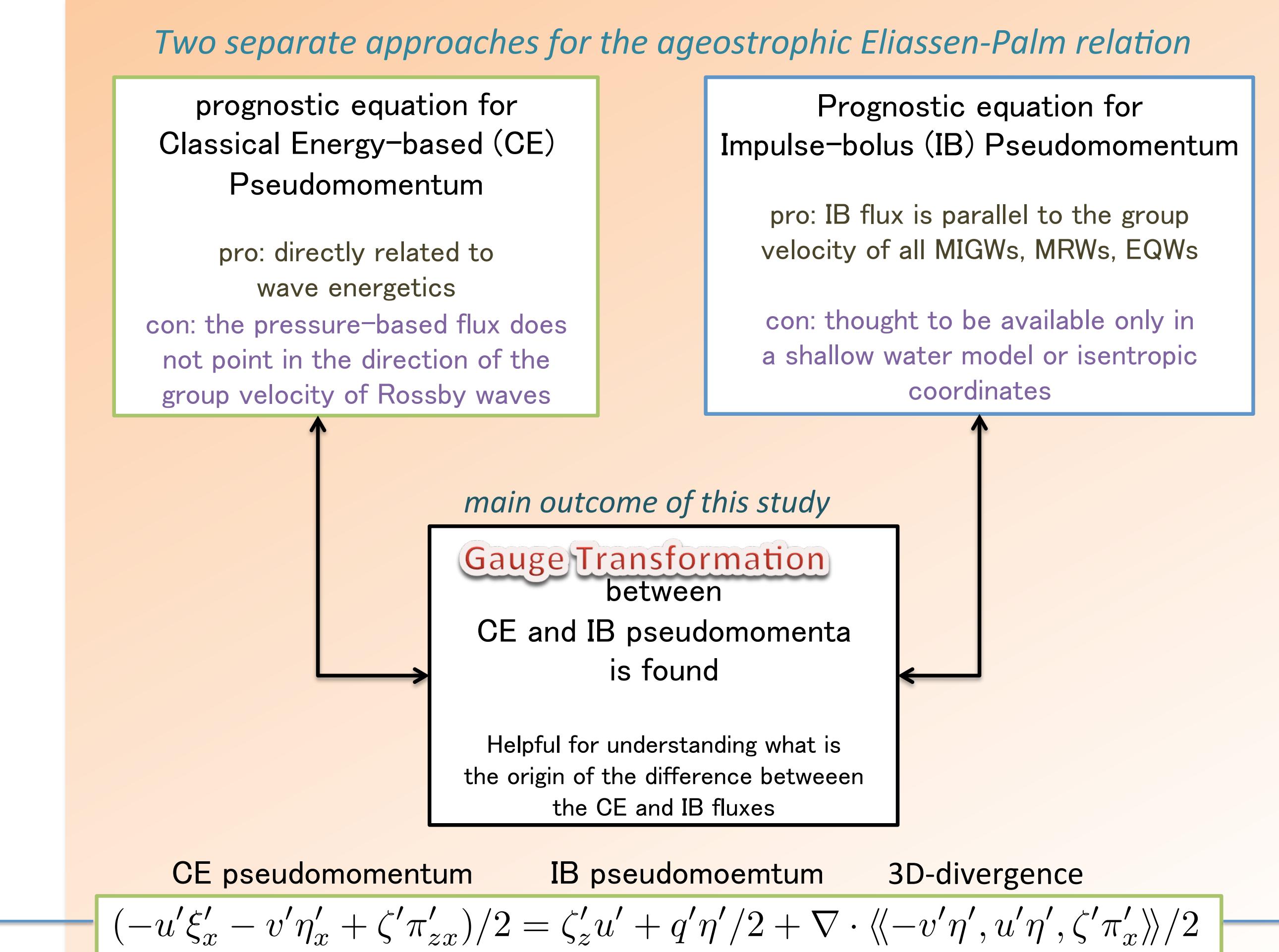
$$\partial_t (\zeta'_z u' + q' \eta'/2) + \nabla \cdot \langle\langle u' u' - K + G, v' u', \zeta' p'_x \rangle\rangle = 0,$$

see gauge transformation

see virial theorem

Bernoulli head in the Lagrangian framework has been thought to be formidable

We suggest to understand it using Lambda in the virial theorem



### Virial Theorem for fluid waves

$$\Lambda \equiv [(\xi' p'_x)_x + (\eta' p'_y)_y + (\zeta' p'_z)_z]/2$$

$$= K + G + f(\xi' v' - \eta' u')/2 - [(\xi' u' - \eta' v')/2]_t$$

$$= -E + (u' u' - \xi' u'_x + v' v'_x - \eta' v'_x)/2 + f(\xi' v' - \eta' u')/2,$$

Second line (left) is the virial theorem represents a well-known equipartition statement between KE and PE associated with linear waves in a nonrotating frame

Lambda has been ignored in almost all previous studies by assuming three-dimensionally homogeneous waves

assuming three-dimensionally homogeneous waves

$\bar{\Lambda} = 0$  for MIGWs: the CE and IB fluxes point in the same direction  
 $\bar{\Lambda} = -E$  for MRWs: the CE and IB fluxes point in different directions

$\bar{\Lambda} \approx \bar{K} - \bar{G} + (f/\beta) \bar{q}' \bar{u}'$   
an expression to be used for model diagnosis in a future study

Previous knowledge for Lambda

New knowledge for Lambda

$$\begin{aligned} \bar{p} + \bar{\Lambda} &= \bar{p} + \xi' p'_x + \eta' p'_y + \zeta' p'_z - \bar{K} \\ &= \bar{p} + \xi' p'_x + \eta' p'_y - \zeta' p'_z - \bar{N}^2 + \bar{G} - \bar{f}(\xi' v' - \eta' u')/2 + [(\xi' u' + \eta' v')/2], \end{aligned}$$

LM pressure Bernoulli head

$\bar{p} + \xi' p'_x + \eta' p'_y - \zeta' p'_z - \bar{N}^2 + \bar{G} - \bar{f}(\xi' v' - \eta' u')/2 + [(\xi' u' + \eta' v')/2]$

LM pressure Bernoulli head

Ageostrophic Taylor-Bretherton identity

$$\begin{aligned} \partial_t (\zeta'_z u') + q' v' &= -\nabla \cdot \langle\langle u' u' - K + G, v' u', \zeta' p'_x \rangle\rangle, \\ \partial_t (\zeta'_z v') - q' u' &= -\nabla \cdot \langle\langle u' v' - K + G, v' v', \zeta' p'_y \rangle\rangle, \end{aligned}$$

Prognostic equation for (what we call) the impulse-bolus (IB) pseudomomentum

$$\begin{aligned} \partial_t (\zeta'_z u' + q' \eta'/2) + \nabla \cdot \langle\langle u' u' - K + G, v' u', \zeta' p'_x \rangle\rangle &= 0, \\ \partial_t (\zeta'_z v' - q' \xi'/2) + \nabla \cdot \langle\langle u' v' - K + G, v' v', \zeta' p'_y \rangle\rangle &= \beta(v' \xi' - u' \eta')/2, \end{aligned}$$

bolus velocity impulse of planetary waves

$$\zeta' (u'_z - w'_x)$$

impulse of gravity waves (nonhydrostatic)

$$\begin{aligned} &\int_{-\infty}^{+\infty} \bar{u}' p'/c_p dy \\ &= \int_{-\infty}^{+\infty} c_g (\bar{E}/c_p) dy \\ &= \int_{-\infty}^{+\infty} c_g (\zeta'_z u' + q' \eta'/2) dy \\ &= \int_{-\infty}^{+\infty} (u' u' - K + G) dy \end{aligned}$$

Merged Eulerian mean (MEM) momentum equations

$$\begin{aligned} \bar{u} + v^{qs} - \zeta'_z u' + q' \eta'/2 + (\xi' v' + \eta' v')_x/2 &+ \nabla \cdot (\bar{U} \bar{u}) - f[\bar{v} + v^{qs} + (\xi' v' - \eta' u')_x/2] &= -(p + \Lambda)_x, \\ \bar{v} + v^{qs} - \zeta'_z v' + q' \xi'/2 + (\xi' u' + \eta' v')_y/2 &+ \nabla \cdot (\bar{U} \bar{v}) + f[\bar{u} + u^{qs} - (\xi' v' - \eta' u')_y/2] &= -(p + \Lambda)_y, \end{aligned}$$

Transformed Lagrangian mean (TLM) momentum equations

$$\begin{aligned} &[\bar{u} + v^{Stokes} - \zeta'_z u' - v' \eta'_x + f(\xi'_z \eta' - \xi' \eta'_z)/2]_t \\ &+ \nabla \cdot (\bar{U} \bar{u}) - f(\bar{v} + v^{Stokes}) - \beta(\eta' v') \\ &= -[(\bar{p} + \xi' p'_x + \eta' p'_y) - \bar{K}] - f(\bar{\xi}' v' - \eta' u')/2, \\ &[\bar{v} + v^{Stokes} - \zeta'_z v' - v' \eta'_y + f(\xi'_z \eta' - \xi' \eta'_z)/2]_t \\ &+ \nabla \cdot (\bar{U} \bar{v}) + f(\bar{u} + u^{Stokes}) + \beta(\eta' u') + \beta(\xi' v' - \eta' u')/2, \\ &= -[(\bar{p} + \xi' p'_x + \eta' p'_y) - \bar{K}] - f(\bar{\xi}' v' - \eta' u')/2, \end{aligned}$$