



# Radiation in the Surface Flux Budget of the NBL: A Question of Timescale

AMS 21<sup>st</sup> BLT Symposium, 8B.5

J. M. Edwards, 10<sup>th</sup> June 2014



# Summary

- Background and motivation
- What determines the **net** LW radiative flux at the surface?
- Idealized model of surface cooling
- Further thoughts

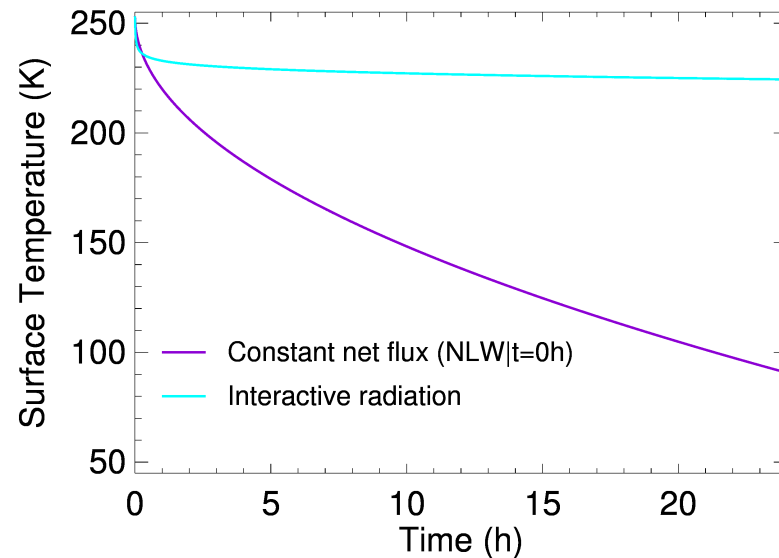


# Background and Motivation

- D. Brunt (1932) “Notes on Radiation in the atmosphere. I.”, QJRMS, vol. 58, p. 389
  - $LW_{\downarrow} = \sigma T^4(a + b\sqrt{e}) \rightarrow R_N = \sigma T^4(1 - a - b\sqrt{e})$
  - *“The vapour pressure has only a slight diurnal variation, and since the fall of temperature during the night is only a relatively small fraction of  $T$ , we may assume as a first approximation that  $R_N$  is constant.”*  
(p. 406)
- If  $R_N$  is constant and  $H=0$ :

$$\Delta T_s = -\frac{2}{\sqrt{\pi}} \frac{R_N}{\sqrt{k_s \rho_s c_s}} \sqrt{t}$$

- Over snow this equation predicts an unrealistically large fall in temperature, so in reality
  - either  $H$  is significant, or (nonexclusively)
  - $R_N$  does decrease
- Interactive radiation model shows fall in  $R_N$  alone can arrest cooling



- GABLS3 SCM intercomparison (Bosveld et al. 2014, BLM)

- Impact of SBL profile on  $LW_{\downarrow}$
- Extension of Brunt's formula

$$L_{ref}^{\downarrow} = (a + b\sqrt{e_{200}})\sigma T_2^4 + (c + d\sqrt{e_{200}})\sigma(T_{200}^4 - T_2^4) + f$$

- For these conditions:
  - Brunt:  $\Delta R_N \approx 0.20 \Delta(\sigma T^4)$
  - Bosveld et al.:  $\Delta R_N \approx 0.26 \Delta(\sigma T^4)$

- A. K. Betts (2006) “Radiative scaling of the nocturnal boundary layer and the diurnal temperature range”, JGR, vol. 111, D07105:
  - *“Surprisingly, comparatively little attention has been paid to the corresponding role of radiative forcing in determining the strength and depth of the NBL...”*
- Betts’ Radiative temperature scale (2004, 2006)
  - $\Delta T_R = -\lambda_0 LW_{n24}$ ;  $\lambda_0 = 1/(4\sigma T^3)$
  - Betts (2006) uses this to scale diurnal temperature range



# The net LW Flux at the surface





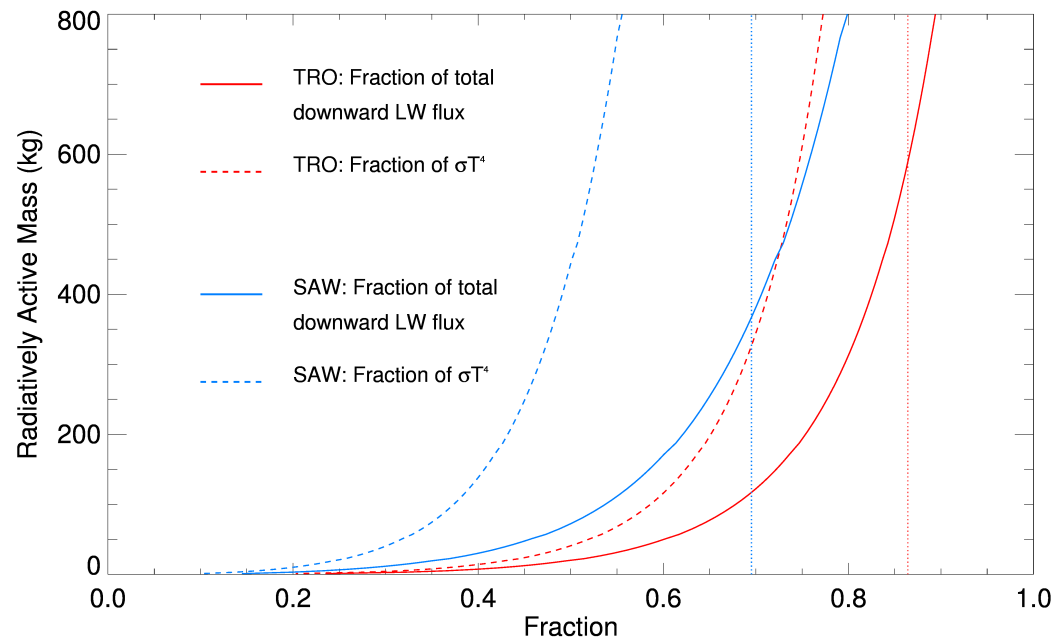
# Downward LW Radiation

- Atmospheric absorbing and emitting gases:
  - Mainly water vapour (lines and continuum)
  - Carbon dioxide
  - Ozone
  - Other trace gases: methane, nitrous oxide, CFCs
- IR range of interest is  $\sim 4 - 1000 \mu\text{m}$ 
  - $8 - 12 \mu\text{m}$  is relatively transparent: “Atmospheric Window”



# Where does LW radiation come from?

- Bosveld et al. (2014) calculate that 50% of LW $\downarrow$  comes from lowest 100m of atmosphere in GABLS3 case
  - See also Ohmura 2001
- Fraction depends on atmosphere

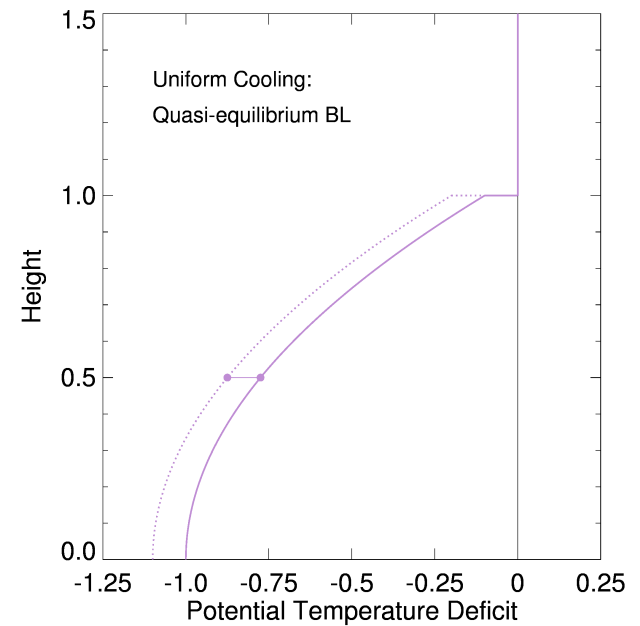
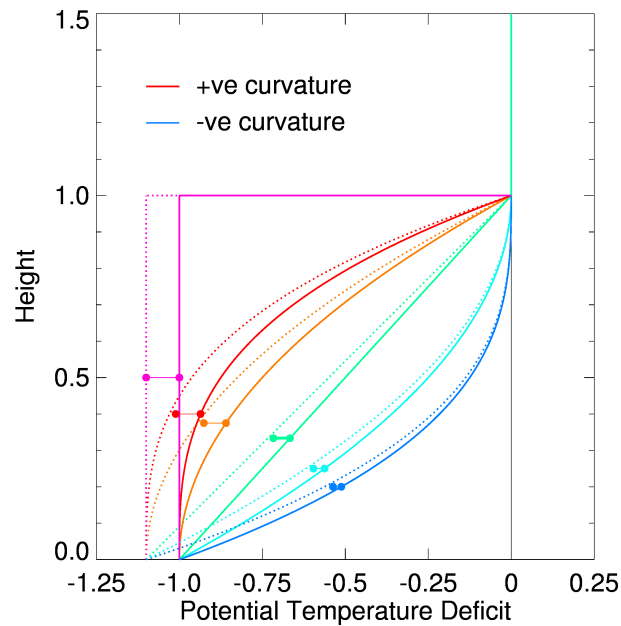


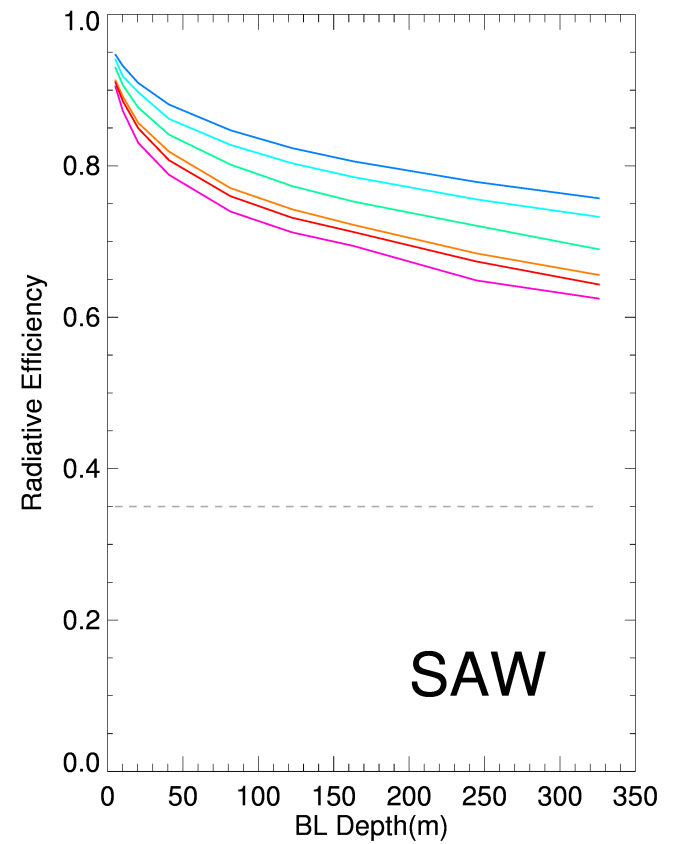
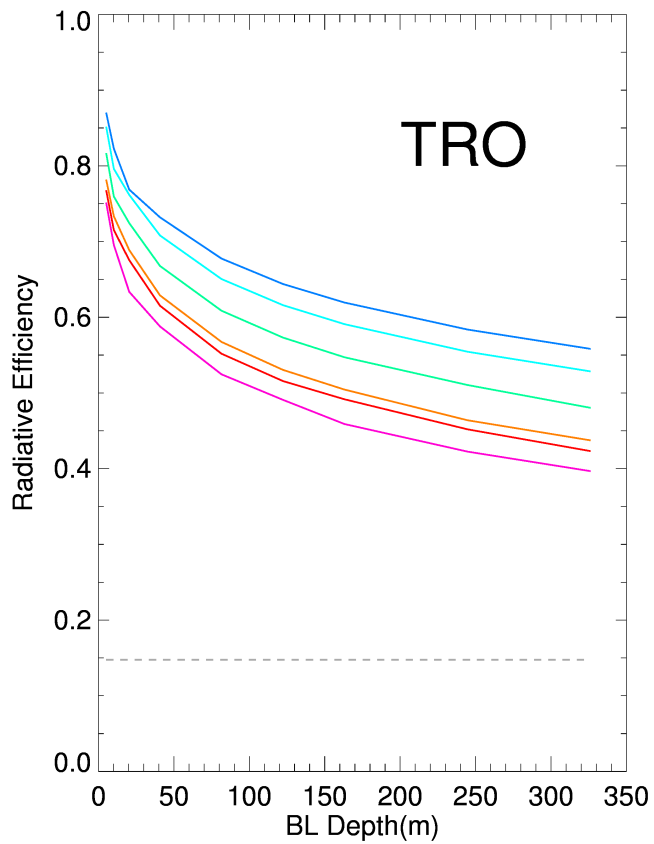


# Radiative Efficiency

- Assume that BL cools in a self-similar way
  - $\Delta\theta(z) = \Delta T(z) = \Delta T_s \text{ fnc}(z)$
- Linearizing net flux,
  - $R_N(t) = R_N(0) + (dR_N/d(\sigma T_s^4)) \Delta(\sigma T_s^4)$
- Define “radiative efficiency”
  - $\eta = dR_N/d(\sigma T_s^4)$
- Using Brunt’s formula,  $\eta = R_N/\sigma T_s^4$ , but not in general

- Now calculate  $\eta$  for different idealized shapes for and depths of the BL







# An idealized model of surface cooling



# Idealized Model of Surface Cooling

- Ground initially at vertically uniform temperature  $T_0$ ; net radiative flux  $R_{NO}$
- Linearize net radiative flux:
  - $R_N = R_{NO} + \eta \Lambda \Delta T_s$ ;  $\Lambda = 4\sigma T_s^3$
  - Defines temperature scale  $R_{NO} / (\eta\Lambda)$  (cf. Betts 2004,2006)
- Surface boundary condition is
  - $k_s dT/dz = R_N = R_{NO} + \eta\Lambda \Delta T_s$ ;  $z$  in soil
  - Implies length-scale  $k_s / (\eta\Lambda)$

- Diffusion in ground,  $\partial T / \partial t = \kappa_s \partial^2 T / \partial z^2$ , implies timescale  $(k_s / (\eta \Lambda))^2 / \kappa_s$
- Non-dimensionalize  $\Delta T$ ,  $z$  and  $t$  with these scales:
  - $\partial \Delta T / \partial t = \partial^2 \Delta T / \partial z^2$
  - $\partial \Delta T / \partial z = 1 + \Delta T; z=0$
  - $\Delta T = 0; t=0$
- Solve by taking Laplace transform



- Solution

$$\Delta T = -1 + \frac{2}{\pi} \int_0^{\infty} \frac{e^{-tu^2}}{1+u^2} du = -1 + e^t (1 - \operatorname{erf}(\sqrt{t}))$$

- Simple approximation ( $\sim 10\%$  error in  $\Delta T$ ):

$$e^{-tu^2} \approx 1 / (1 + tu^2)$$

- Giving

$$\Delta T \approx -\sqrt{t} / (1 + \sqrt{t})$$

- In dimensional units

$$\Delta T \approx \frac{-N_0}{\sqrt{\kappa\rho c}} \sqrt{t} \left[ 1 + \sqrt{t/t_D} \right]^{-1}$$

- Radiation plays a significant role in arresting fall in  $T_s$  if lifetime of SBL is long compared to  $t_D$ , ie.

$$t_{SBL} > t_D = \frac{\rho_s c_s k_s}{\eta^2 \Lambda^2}$$

- Some typical values of  $t_D$ 
  - Deep humid tropical NBL:  $t_D \approx 10$  days
  - Subtropical desert NBL:  $t_D \approx 7$  hours
  - Polar SBL over fresh snow:  $t_D \approx 1$  hour

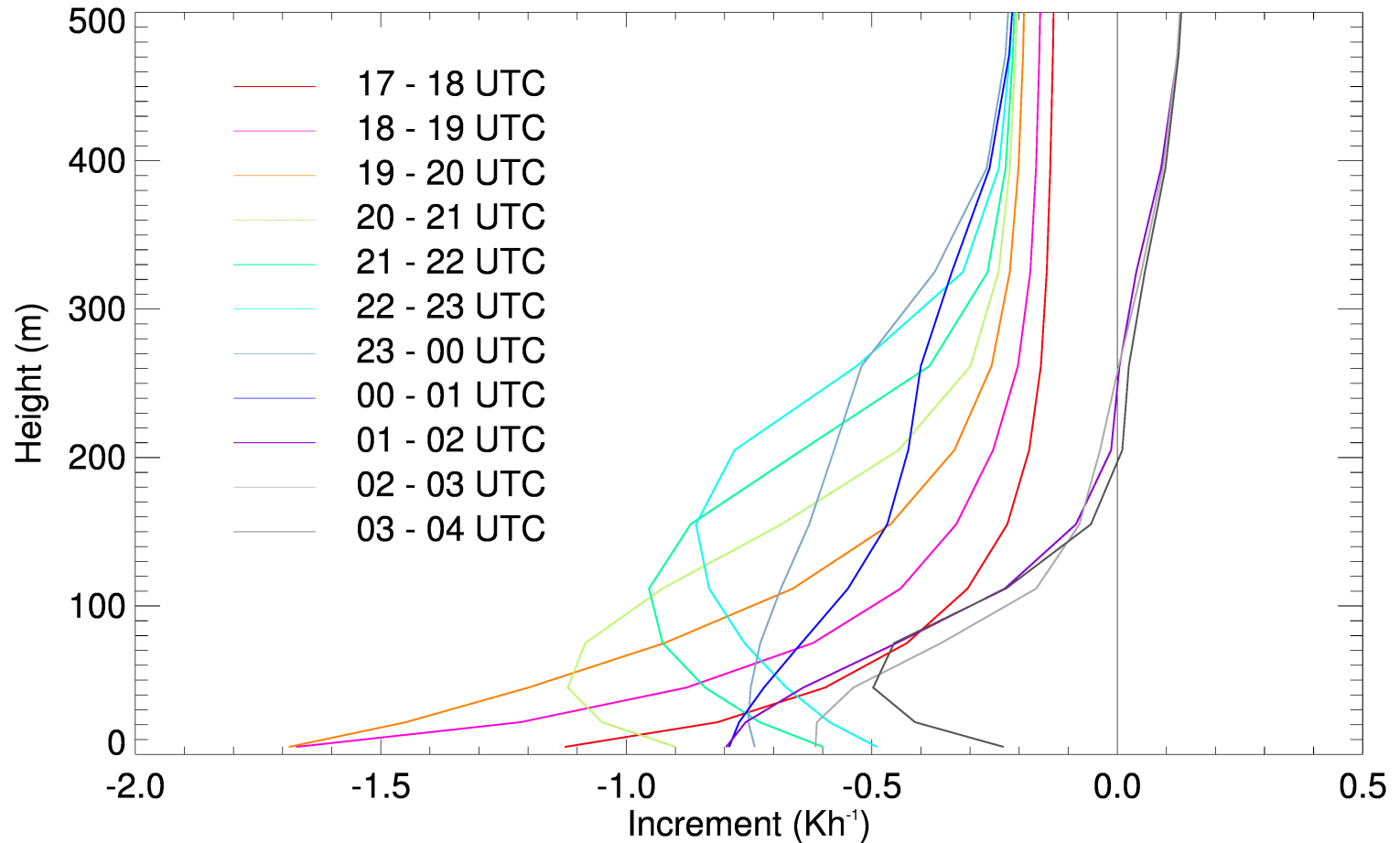


# Further thoughts



- This is an idealized model
  - More illustrative than predictive
- Caveats:
  - Cooling of the residual layer is not included, or assumed to be compensated by subsidence
  - $\eta$  is at root a measure of the ratio of the BL depth to the photon mean free path -- but appropriately weighted in frequency!
  - Actual SBLs are not exactly self-similar, so the efficiency is not rigorously defined

# Hourly temperature increments in GABLS3





# Questions and answers