



Strong Winds in Extratropical Cyclones: Sting Jets and Cold Conveyor Belts

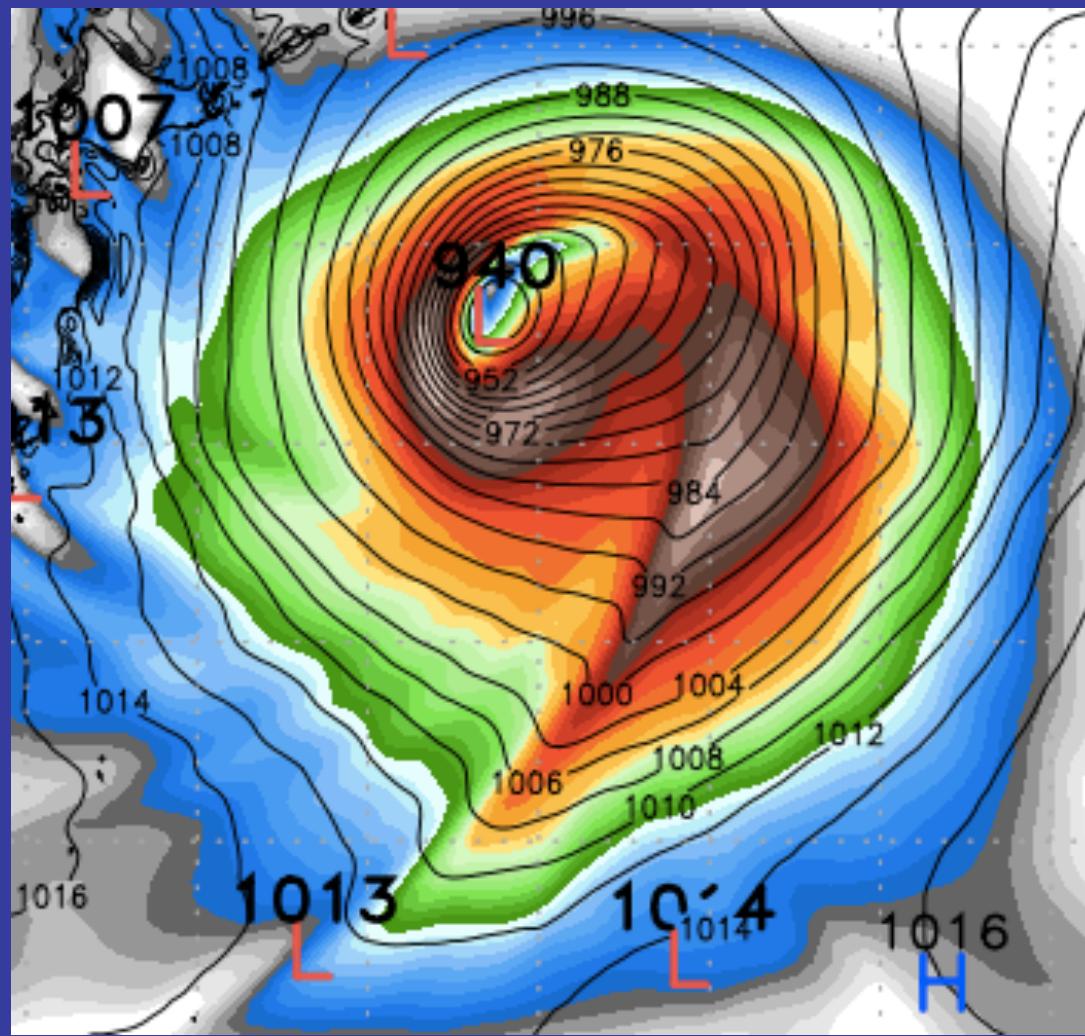
David M. Schultz

University of Manchester

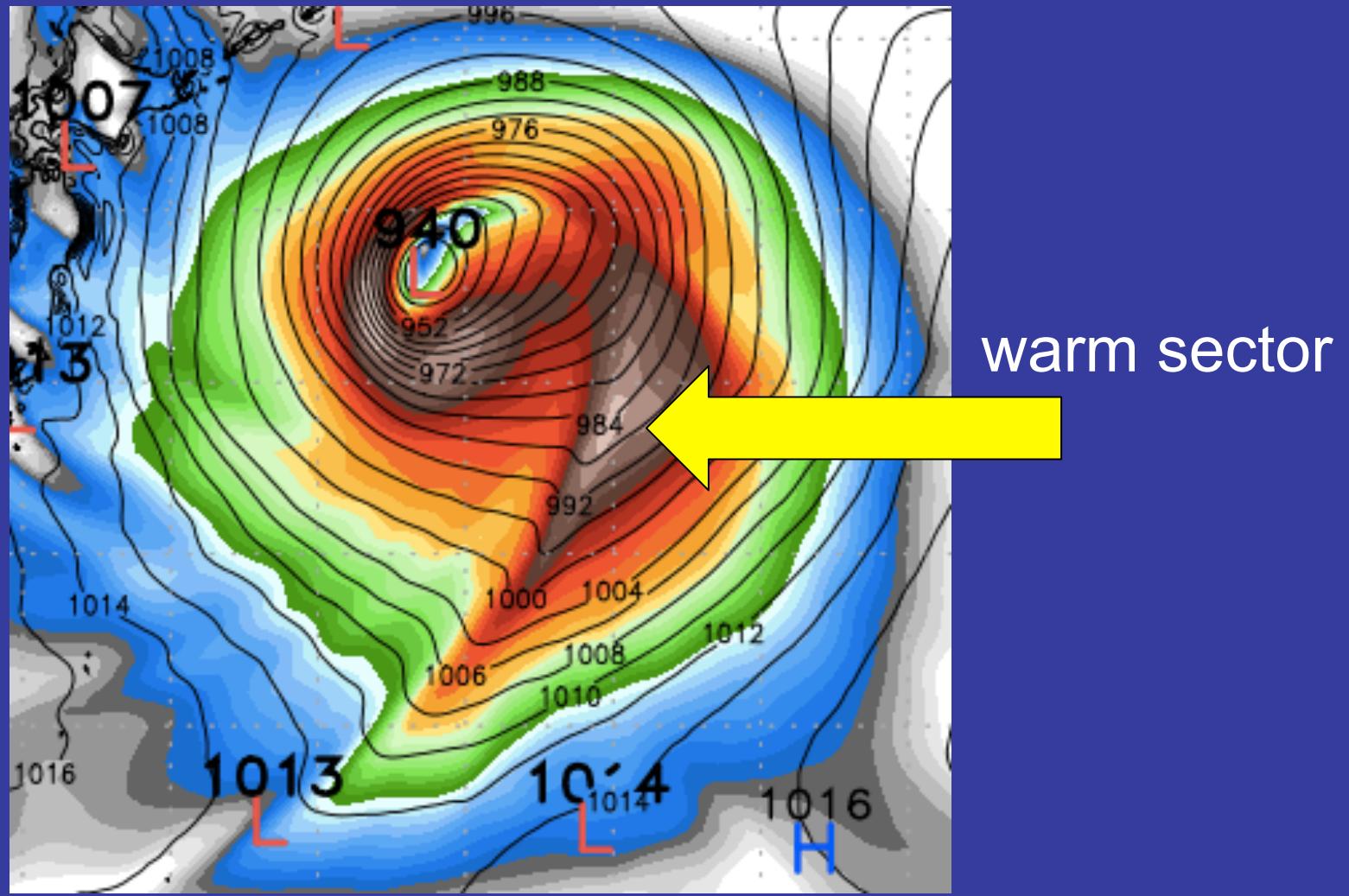
@EloquentScience

with J. Sienkiewicz, G. Vaughan, T. Slater

Three locations where strong winds occur in extratropical cyclones.

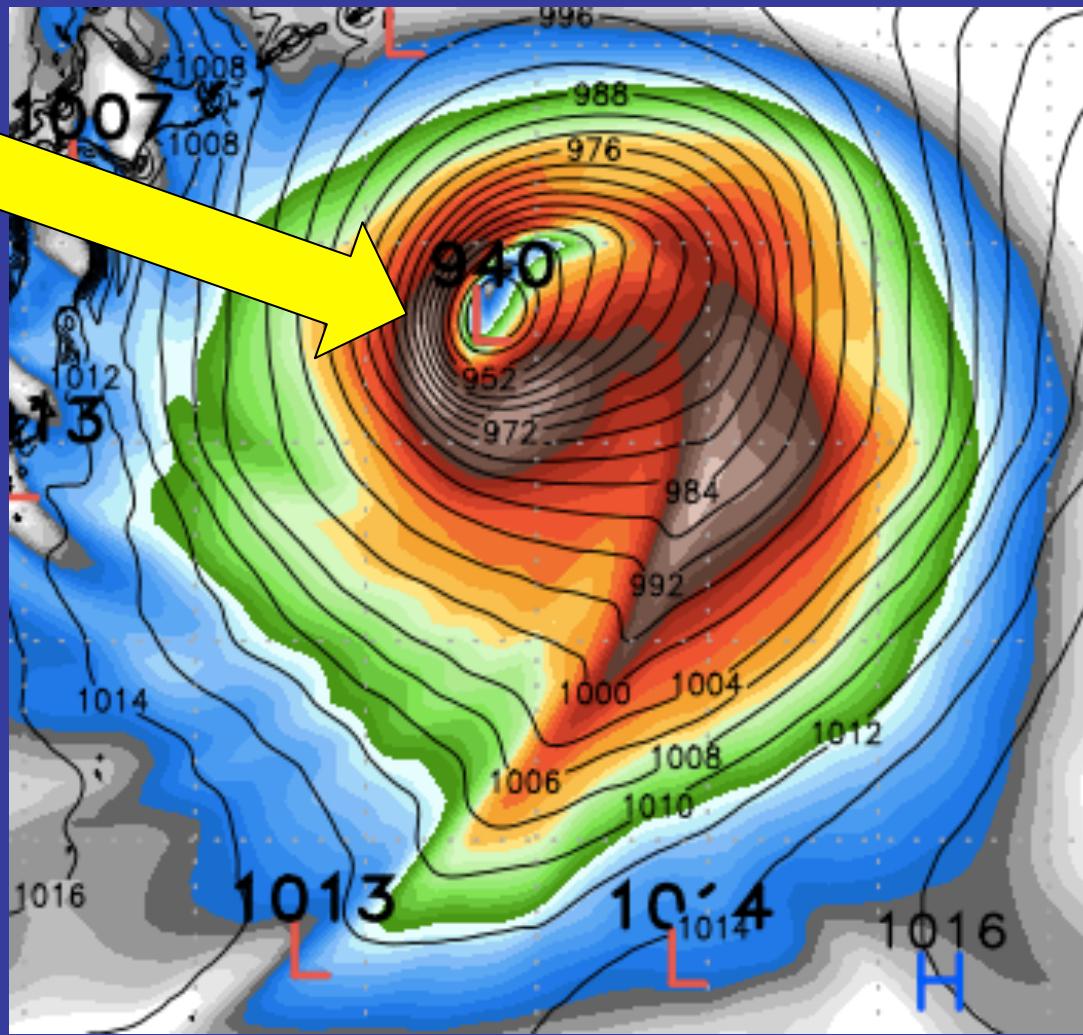


Three locations where strong winds occur in extratropical cyclones.

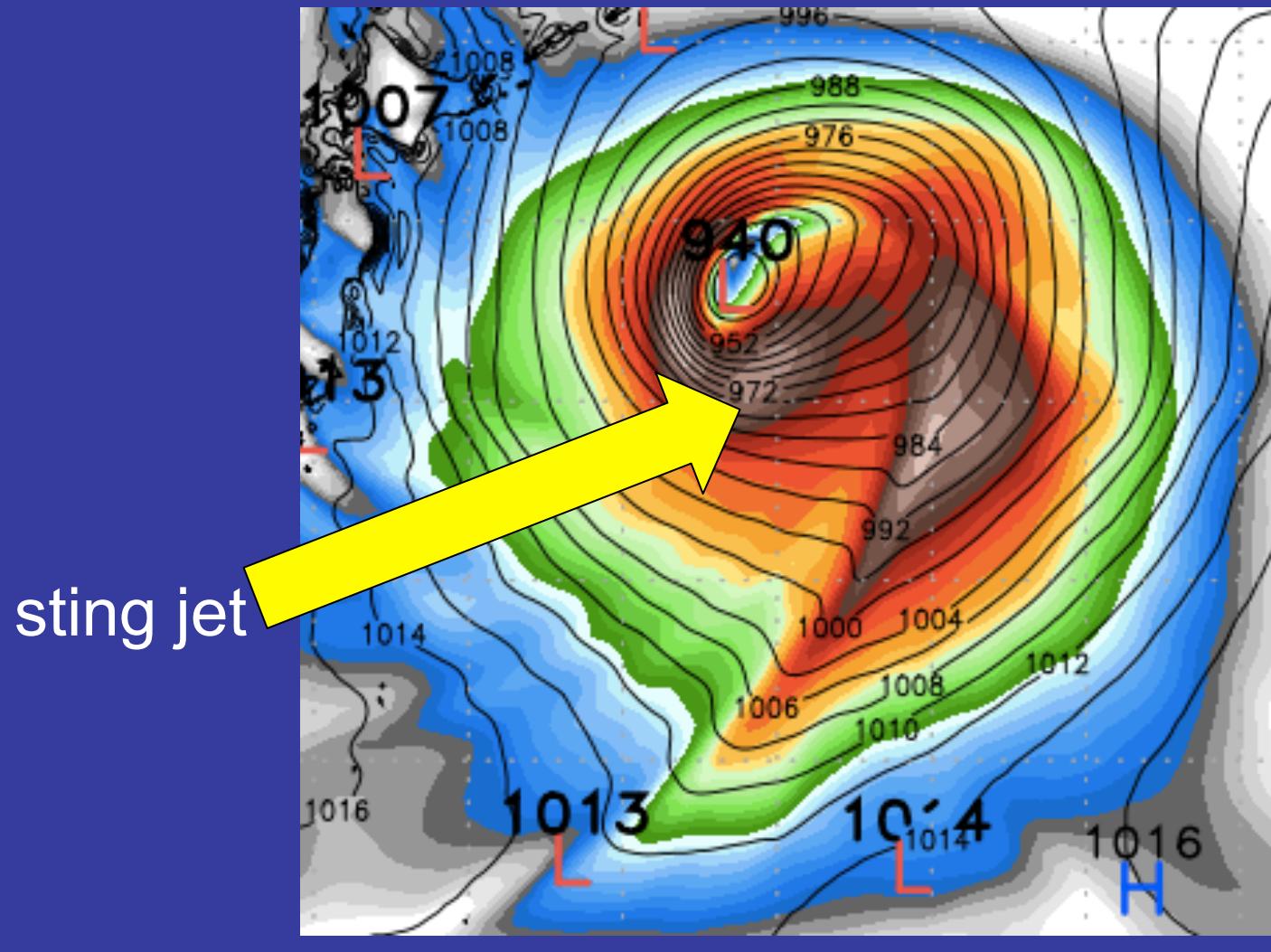


Three locations where strong winds occur in extratropical cyclones.

cold
conveyor
belt

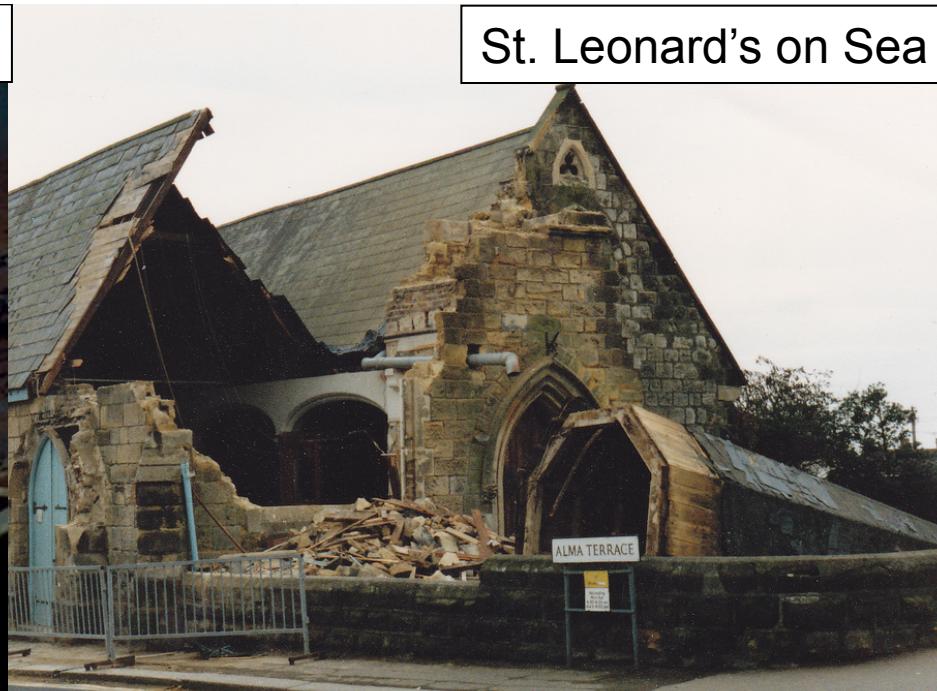


Three locations where strong winds occur in extratropical cyclones.





Folkestone



St. Leonard's on Sea



Peacehaven

16 October 1987

BBC

TONIGHT



Sting jet technology means no more hurricane mishaps for Michael Fish



Follow

Storm #Frank is developing a sting jet (think 1987 storm), this thankfully staying well out to sea. Chris F



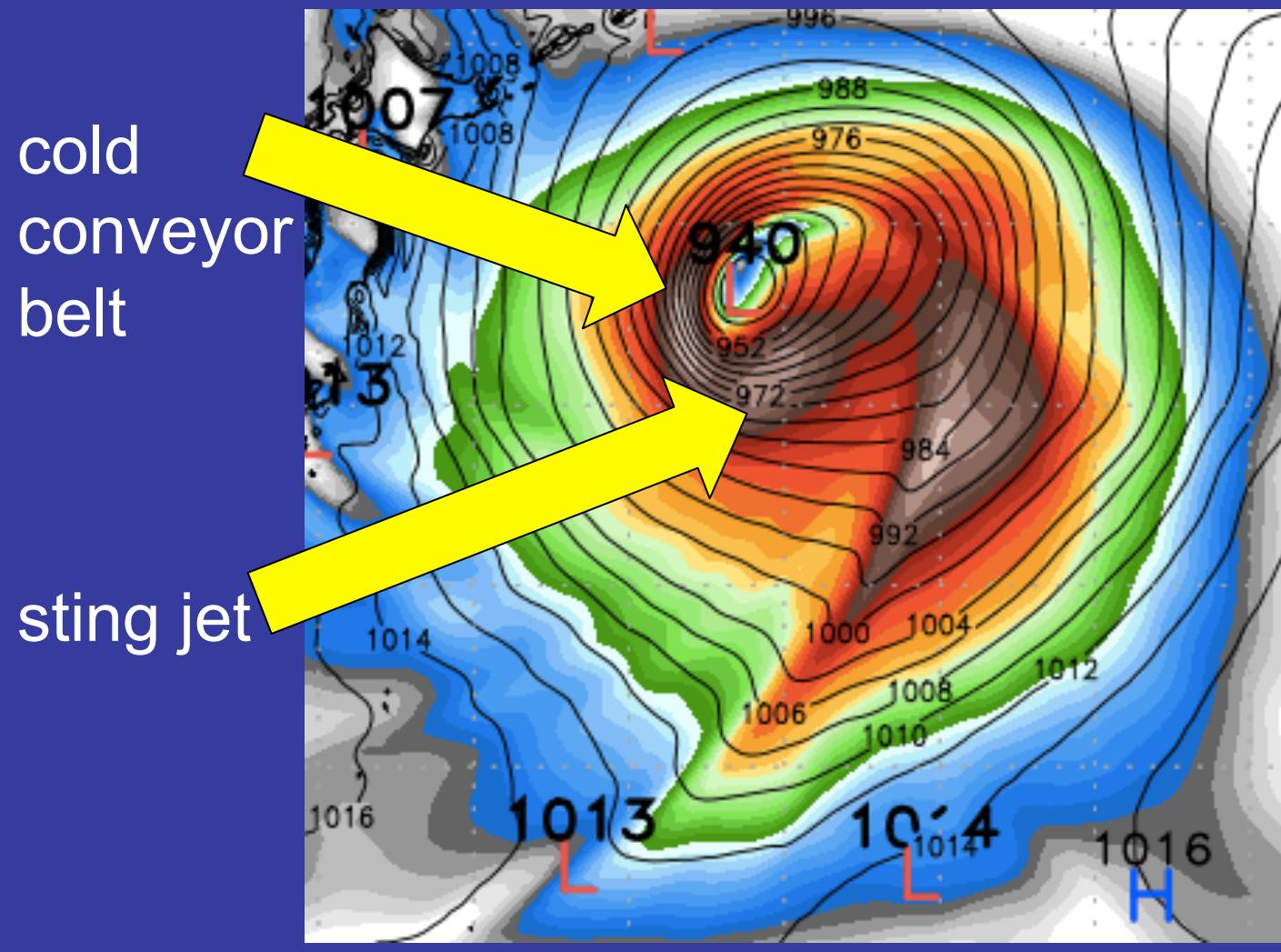
the guardian



risk for power outages early tomorrow morning. The part of the storm we have to pay particular close attention to houses the highest pressure gradients—or the back-bent occlusion/ trough that wraps itself around the low pressure system. In rare cases, this can form the 'poisonous tail of the back-bent occlusion,' or also simply referred to as a sting jet.

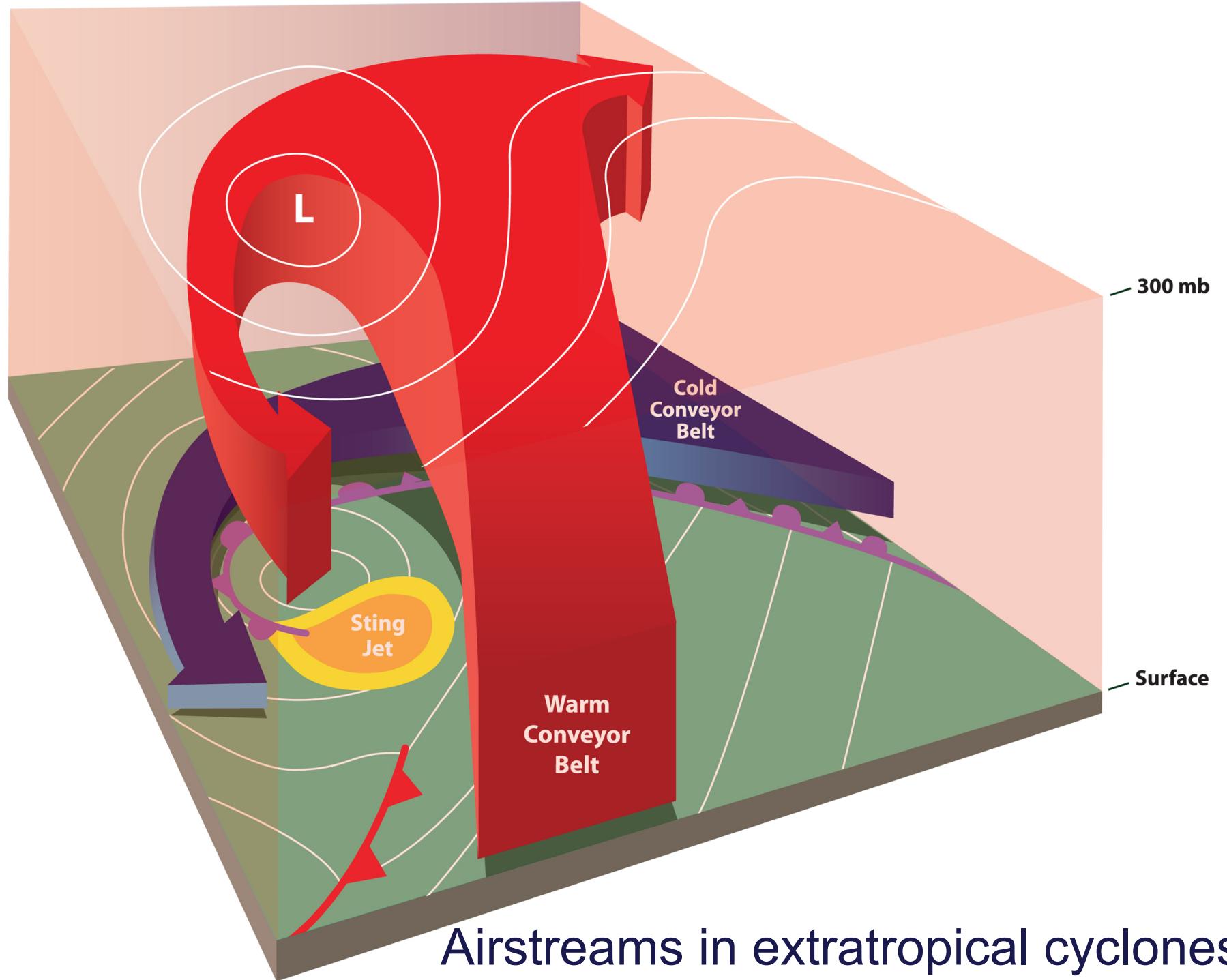


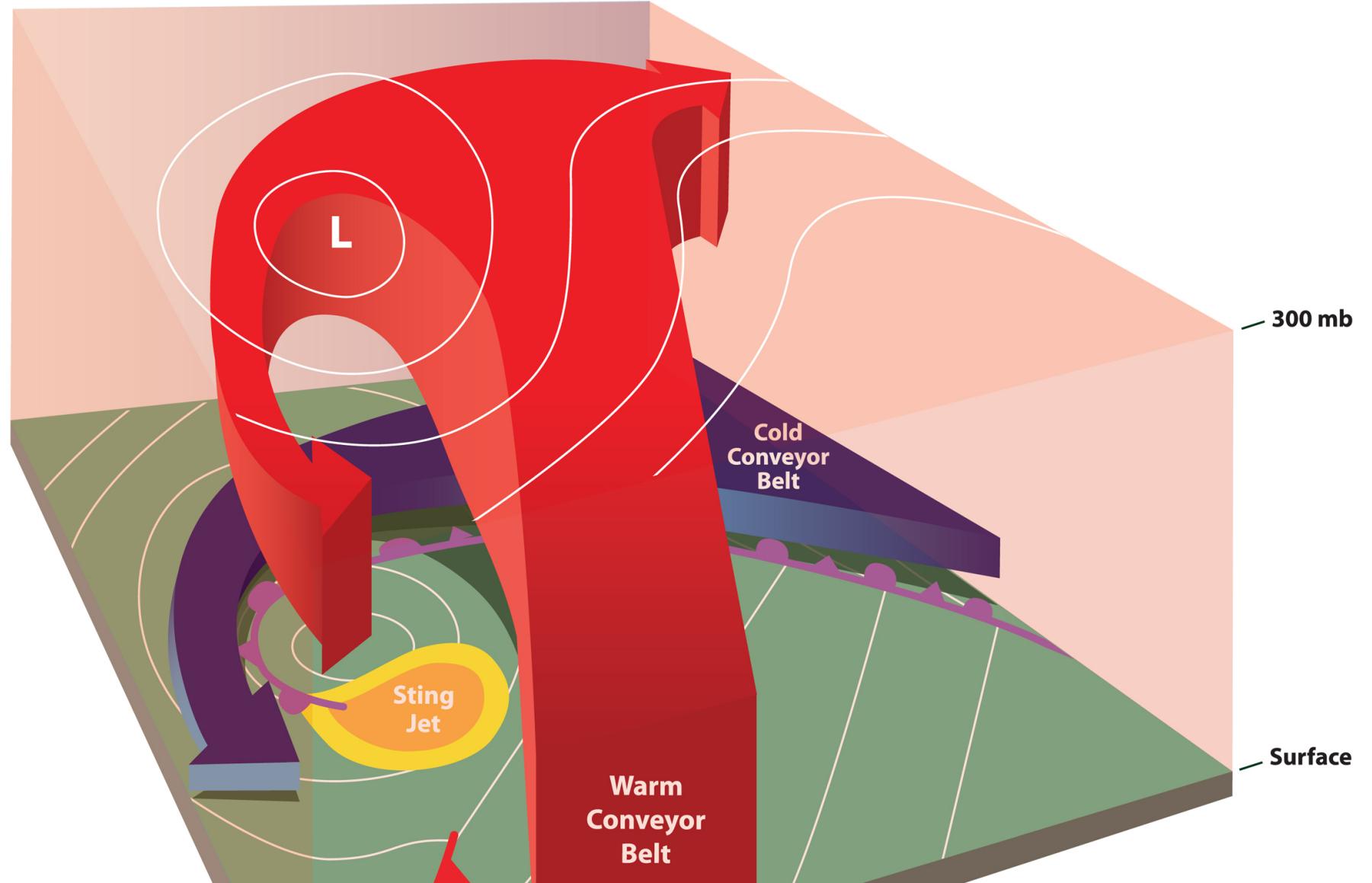
Clear spatial distinction, but not always.



*What distinguishes the sting jet
from the cold conveyor belt?*

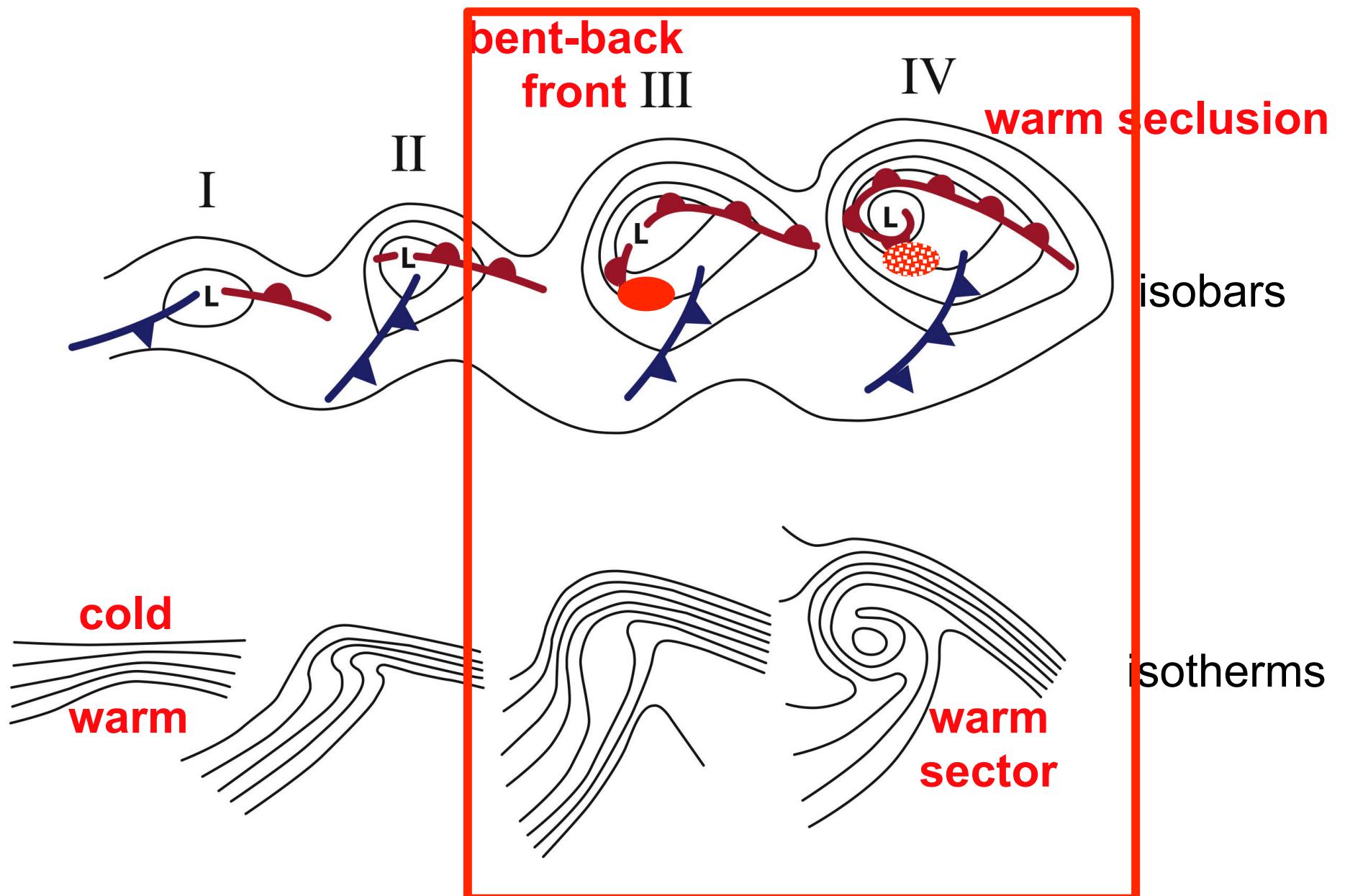
(Schultz and Browning 2017, *Weather*)

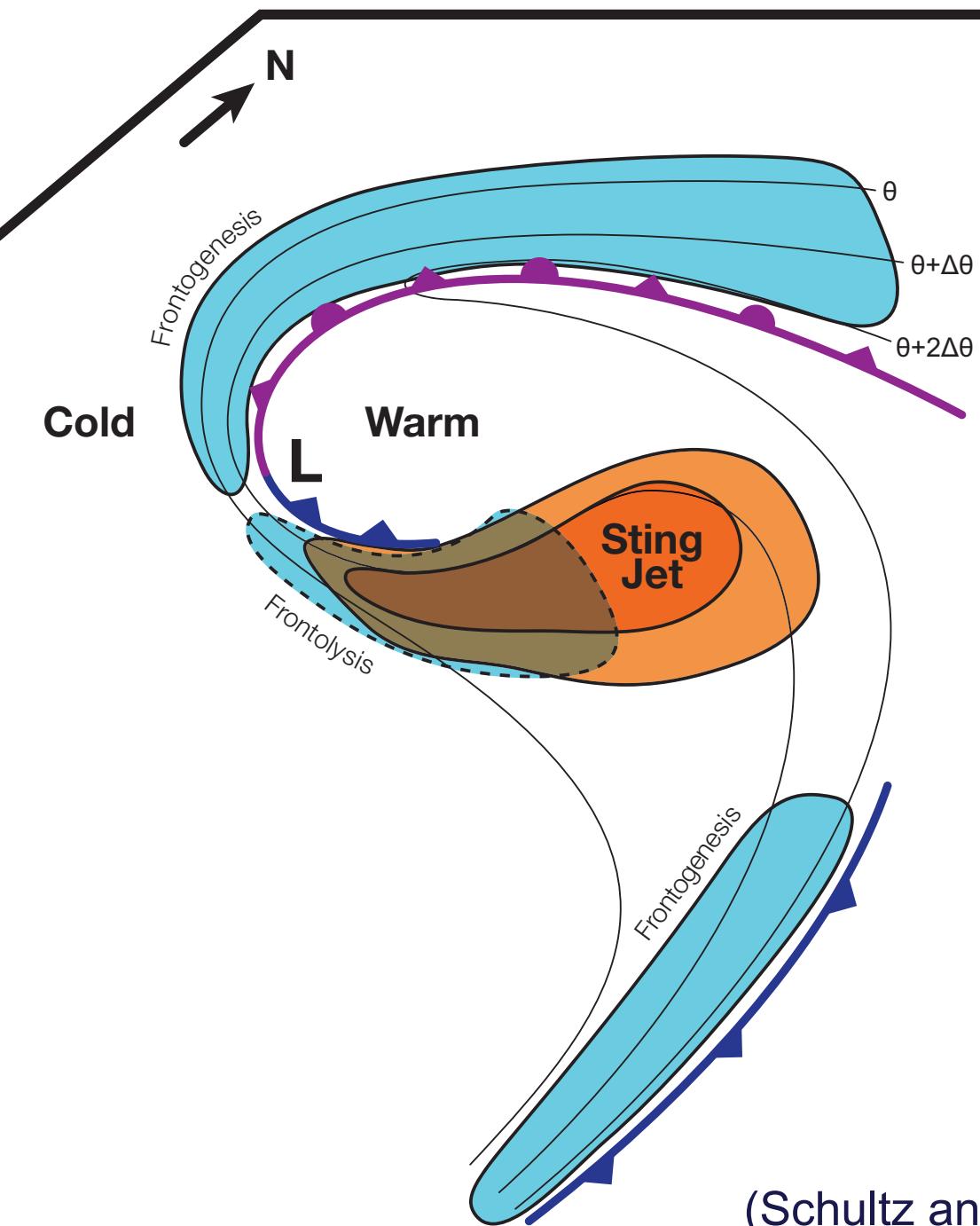




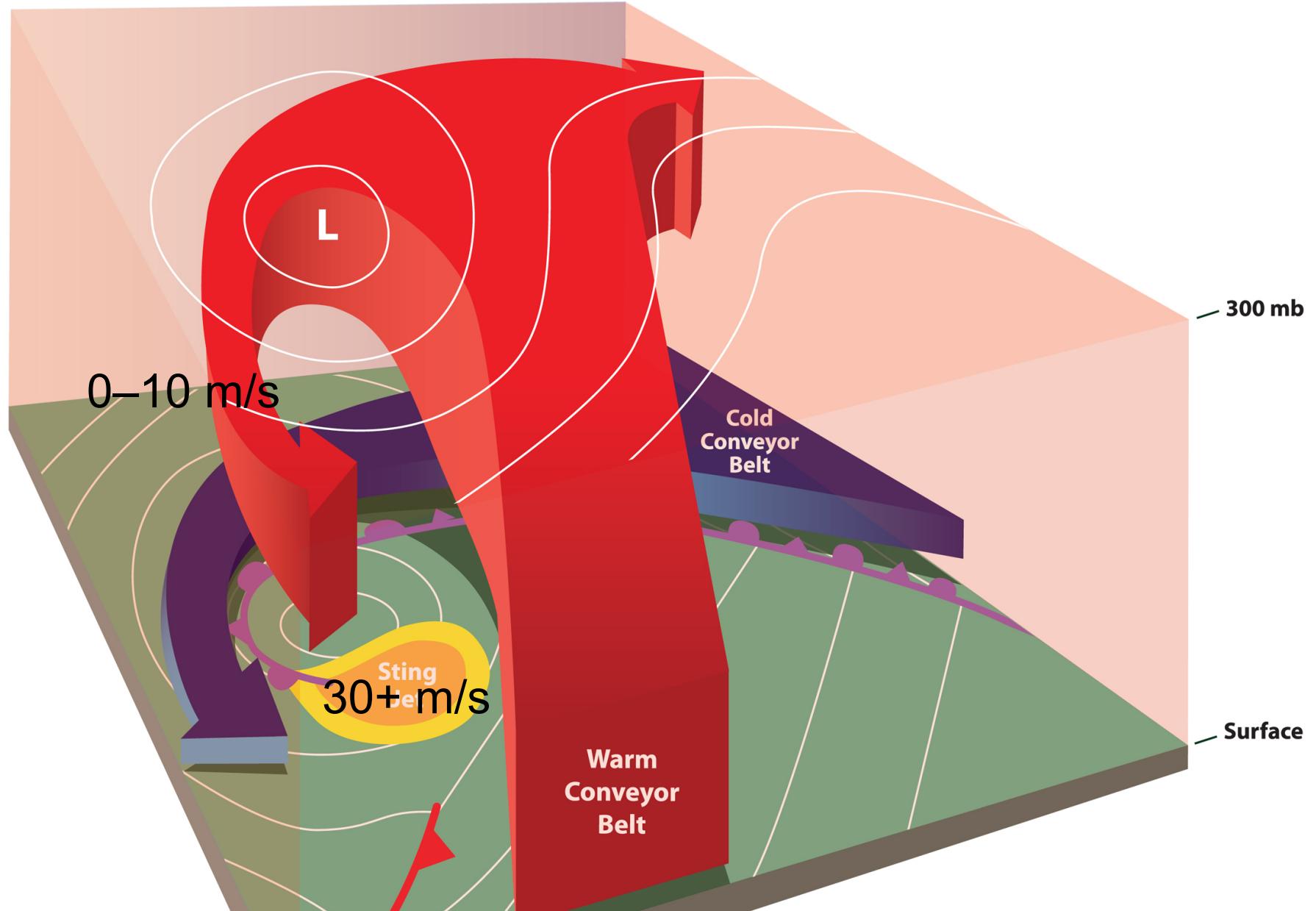
The sting jet descends and occurs at the tail of the bent-back front.

Sting jet in context of Shapiro–Keyser (1990) cyclones





(Schultz and Sienkiewicz 2013)

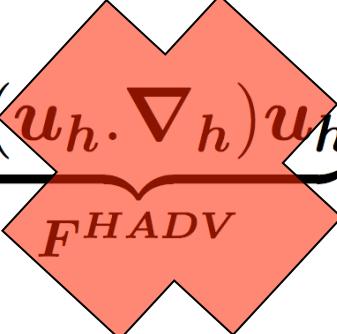


Why does the air accelerate during
downslope?

The momentum equation tells us how winds accelerate.

$$\frac{\partial \mathbf{u}_h}{\partial t} = \underbrace{-(\mathbf{u}_h \cdot \nabla_h) \mathbf{u}_h}_{F^{HADV}} - w \cdot \underbrace{\frac{\partial \mathbf{u}_h}{\partial z}}_{F^{VADV}} - f \underbrace{\mathbf{k} \times \mathbf{u}_h}_{F^{COR}} - \underbrace{\frac{1}{\rho} \nabla_h p}_{F^{PGF}} + \mathbf{F}^{PBL}$$

The momentum equation tells us how winds accelerate.

$$\frac{\partial \mathbf{u}_h}{\partial t} = -(\mathbf{u}_h \cdot \nabla_h) \mathbf{u}_h - w \cdot \frac{\partial \mathbf{u}_h}{\partial z} - f \mathbf{k} \times \mathbf{u}_h - \frac{1}{\rho} \nabla_h p + \mathbf{F}^{PBL}$$


The diagram illustrates the momentum equation components. The first term, $(\mathbf{u}_h \cdot \nabla_h) \mathbf{u}_h$, is highlighted with a red diamond and labeled F^{HADV} . The other terms are grouped under F^{VADV} , F^{COR} , and F^{PGF} .

just moves momentum around:
can't explain acceleration

The momentum equation tells us how winds accelerate.

$$\frac{\partial \mathbf{u}_h}{\partial t} = -(\mathbf{u}_h \cdot \nabla_h) \mathbf{u}_h - w \cdot \frac{\partial \mathbf{u}_h}{\partial z} - f \mathbf{k} \times \mathbf{u}_h - \frac{1}{\rho} \nabla_h p + \mathbf{F}^{PBL}$$

The diagram illustrates the momentum equation components. It shows four red diamond-shaped terms representing different forces: F^{HADV} , F^{VADV} , F^{COR} , and F^{PGF} . These terms are grouped by black brackets below them: F^{HADV} and F^{VADV} are grouped together under a bracket labeled F^{HADV} ; F^{COR} and F^{PGF} are grouped together under a bracket labeled F^{PGF} . The bracket for F^{HADV} also covers the term $w \cdot \frac{\partial \mathbf{u}_h}{\partial z}$.

acts perpendicular to winds:
can't explain acceleration

The momentum equation tells us how winds accelerate.

$$\frac{\partial \mathbf{u}_h}{\partial t} = -(\mathbf{u}_h \cdot \nabla_h) \mathbf{u}_h - w \cdot \frac{\partial \mathbf{u}_h}{\partial z} - f k \times \mathbf{u}_h - \frac{1}{\rho} \nabla_h p + \mathbf{F}^{PBL}$$

The diagram illustrates the momentum equation components. It shows four red diamond-shaped terms: F^{HADV} , F^{VADV} , F^{COR} , and F^{PGF} . These terms are grouped by brackets below them: F^{HADV} and F^{VADV} are grouped together under a bracket labeled F^{HADV} ; F^{COR} and F^{PGF} are grouped together under a bracket labeled F^{PGF} .

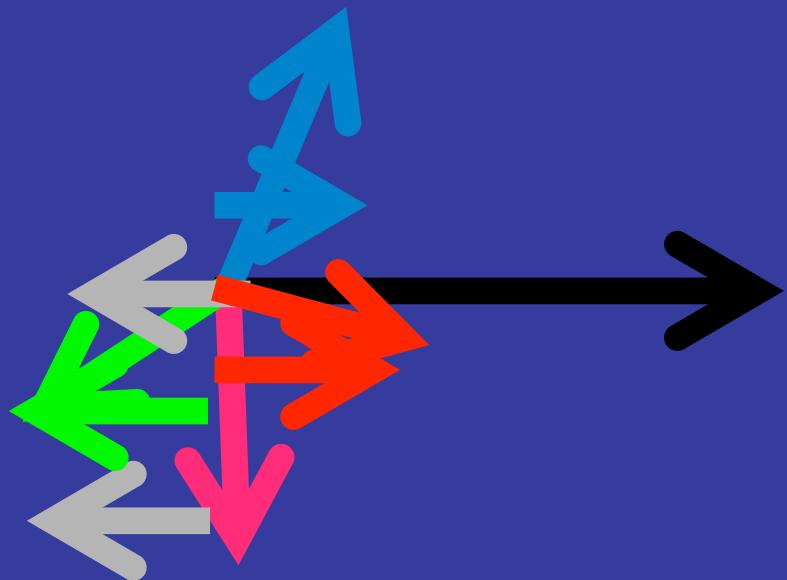
slows wind down:
can't explain acceleration

The momentum equation tells us how winds accelerate.

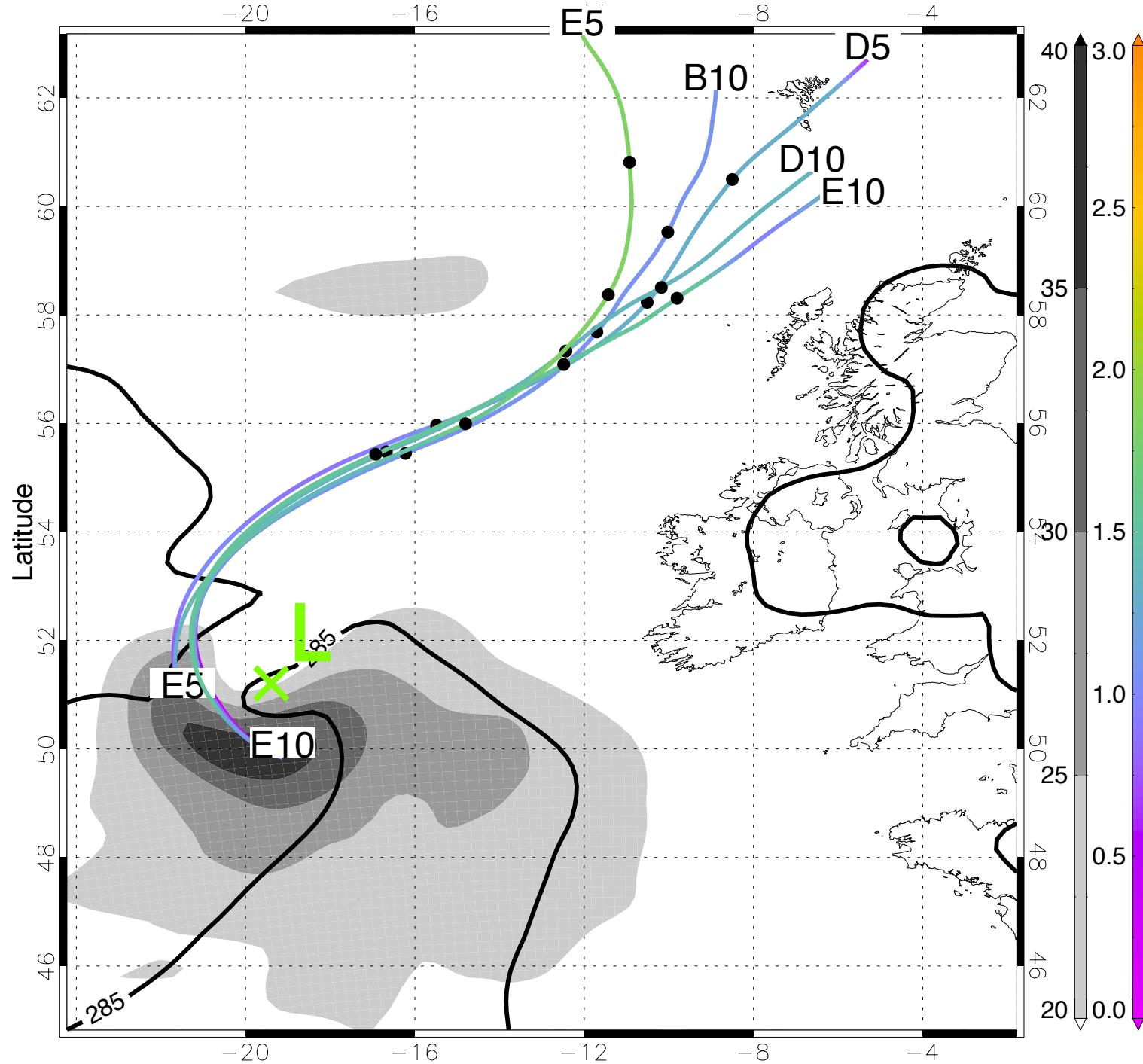
$$\frac{\partial \mathbf{u}_h}{\partial t} = -(\mathbf{u}_h \cdot \nabla_h) \mathbf{u}_h + \underbrace{-w \cdot \frac{\partial \mathbf{u}_h}{\partial z}}_{\mathbf{F}^{VADV}} + \underbrace{-f \mathbf{k} \times \mathbf{u}_h}_{\mathbf{F}^{COR}} + \underbrace{-\frac{1}{\rho} \nabla_h p}_{\mathbf{F}^{PGF}} + \mathbf{F}^{PBL}$$

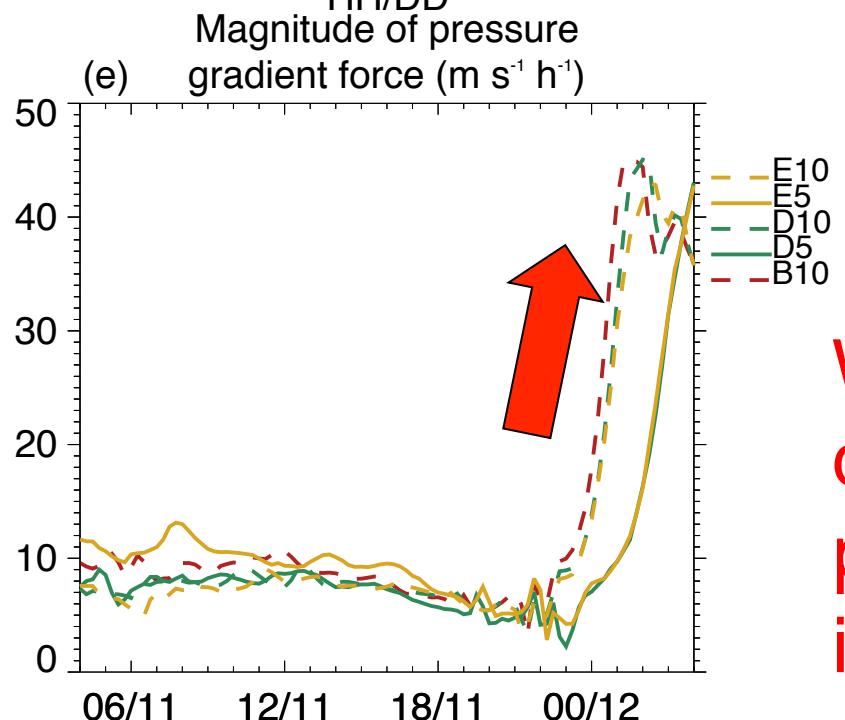
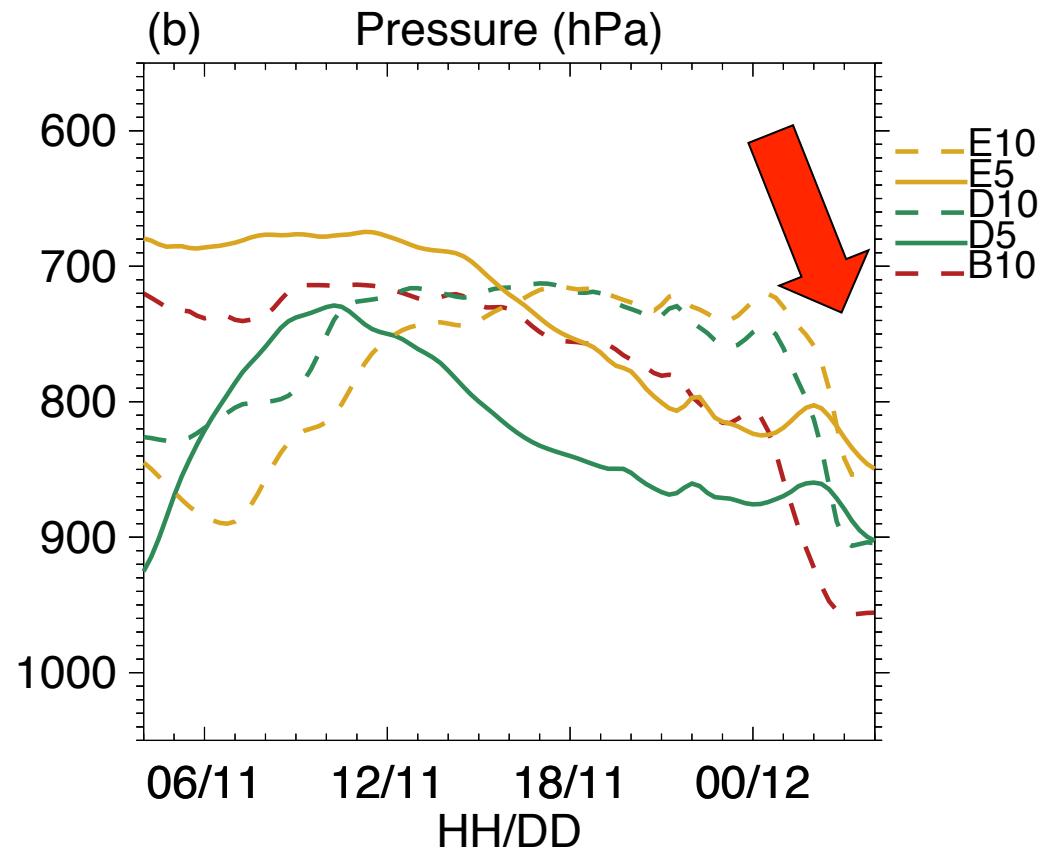
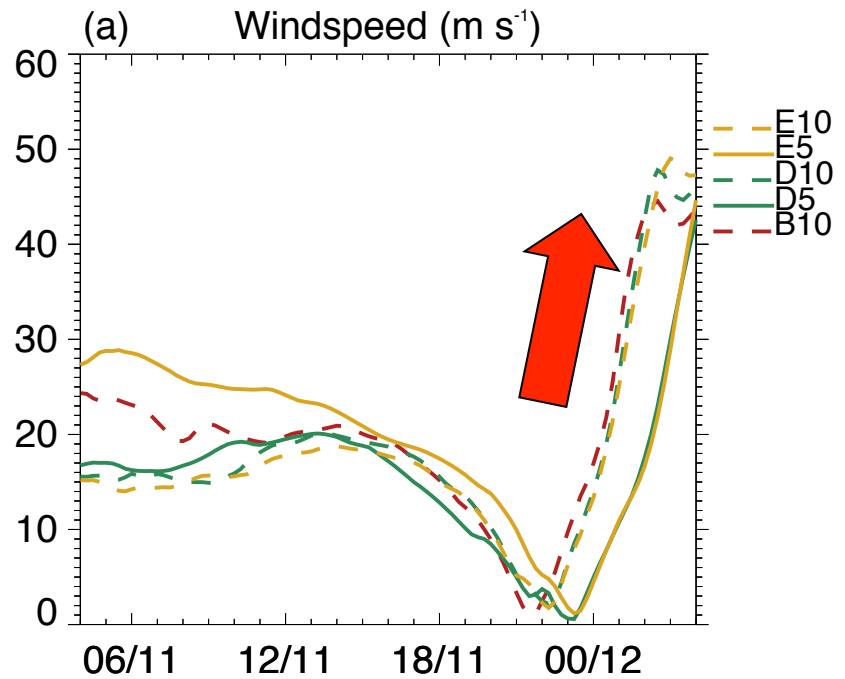
Can create acceleration/force vectors from model output.

To understand increases in wind speed look at along-flow accelerations.

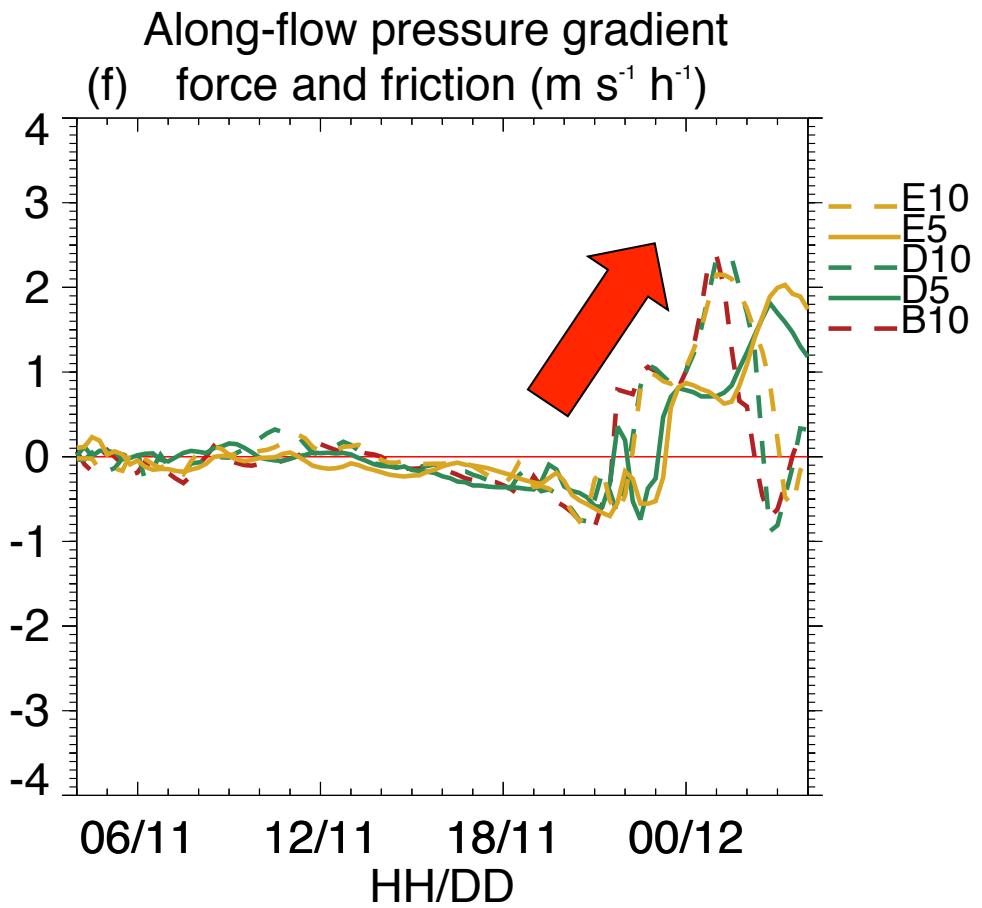
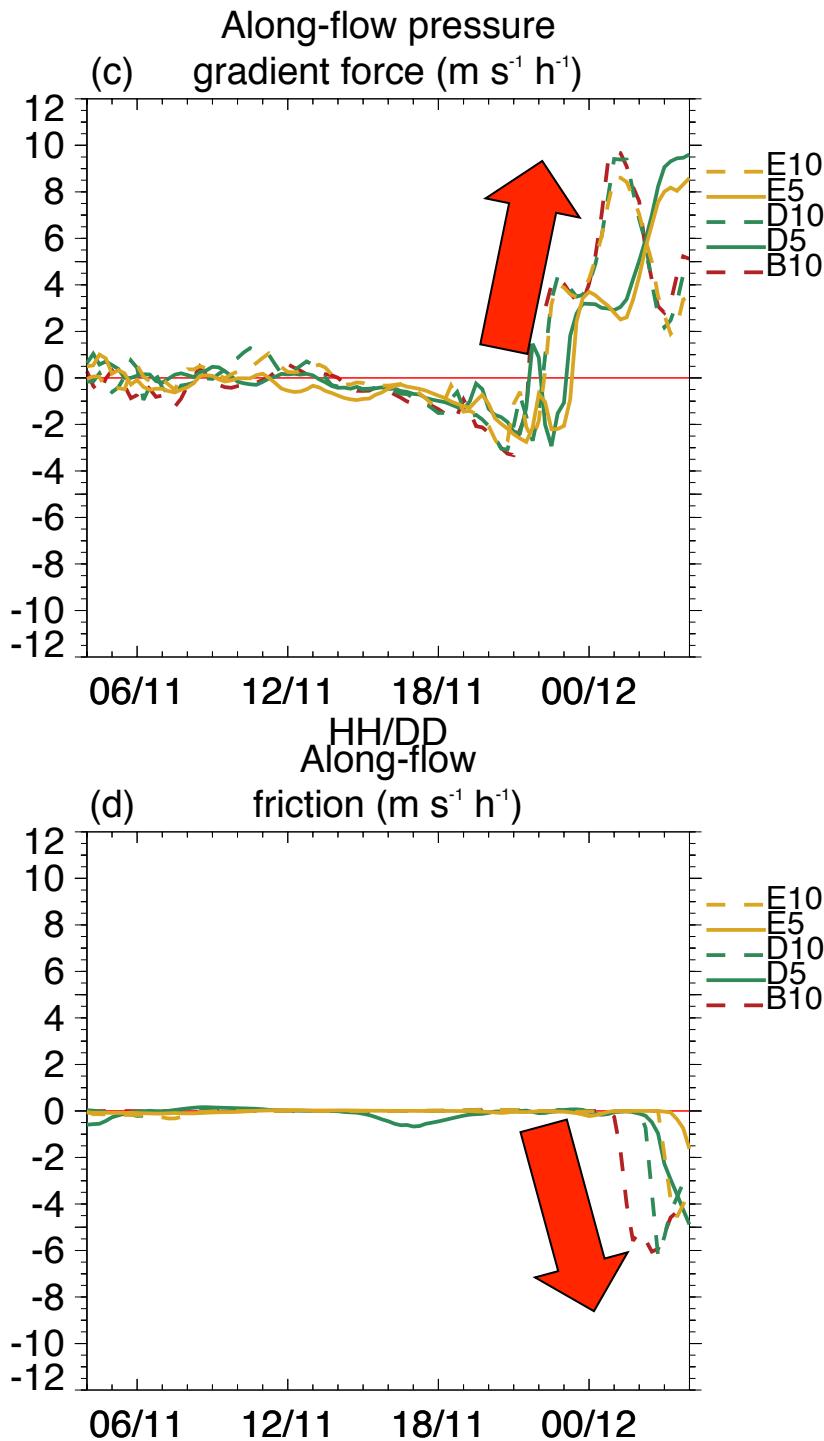


- Velocity
- Pressure gradient force
- Coriolis force
- Vertical advection
- Horizontal advection
- Friction





Wind speed increases during descent as the horizontal pressure gradient force increases substantially.



Acceleration largely happens aloft, but near-surface friction is not enough to decelerate descending air.

Strong winds in cyclones

Frontolysis associated with descent of the sting jet.

- End of bent-back front
- Lasts for a finite time just before maturity of cyclone

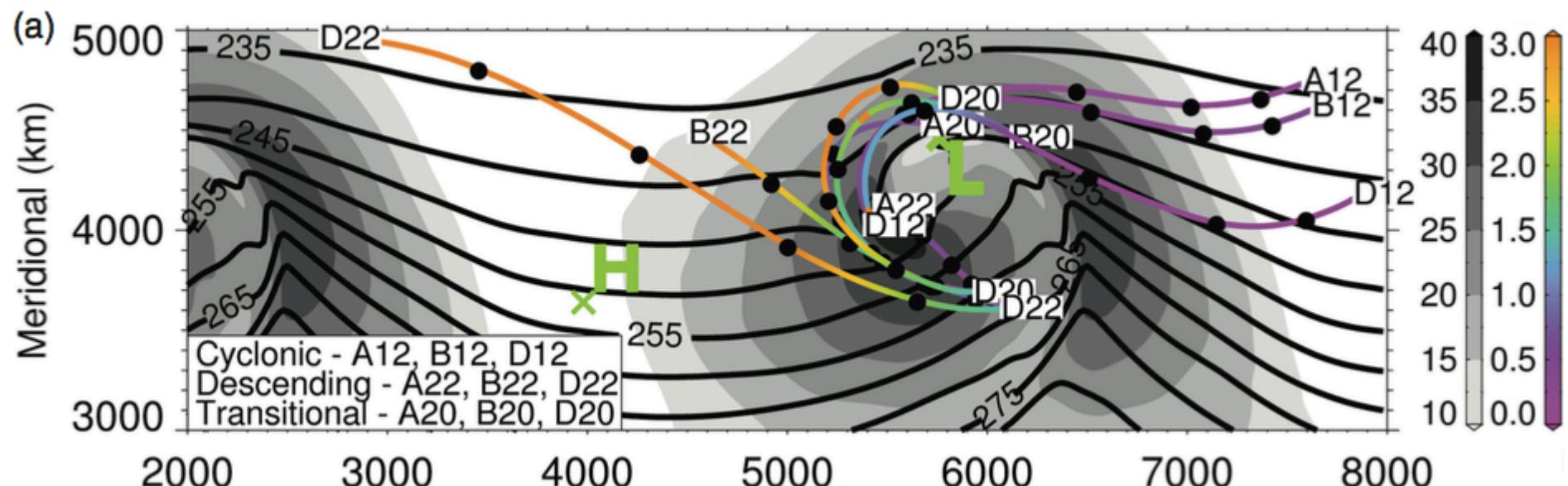
Southwest wind maximum accelerated by both:

- Horizontal pressure-gradient force
- Downward advection of momentum.

Complicates identifying “sting jet” in real time operations.

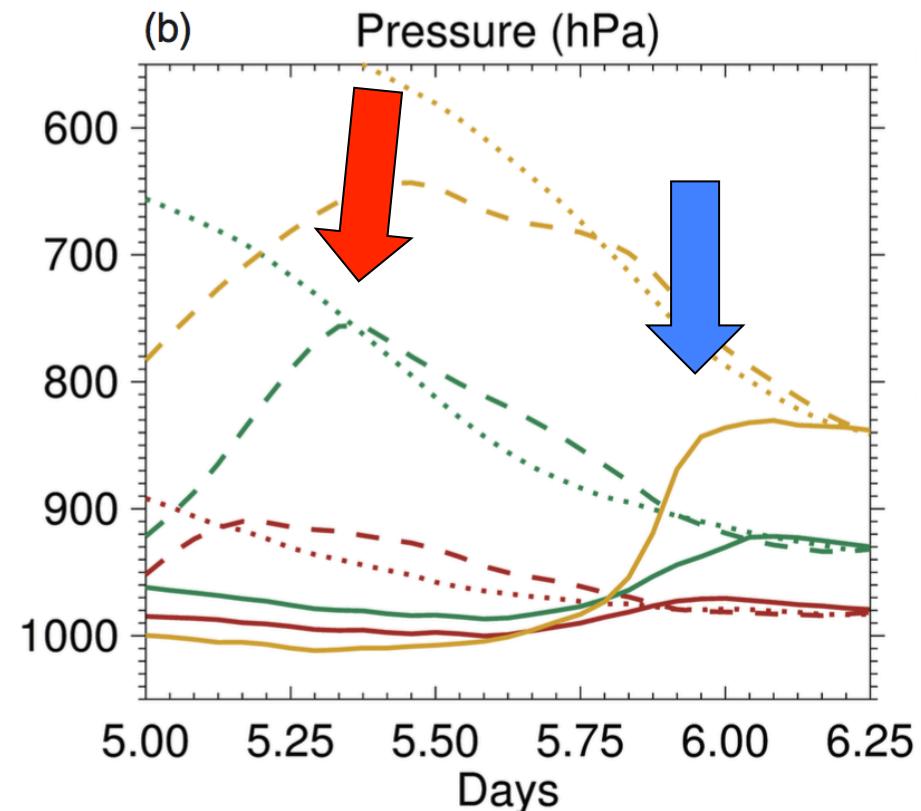
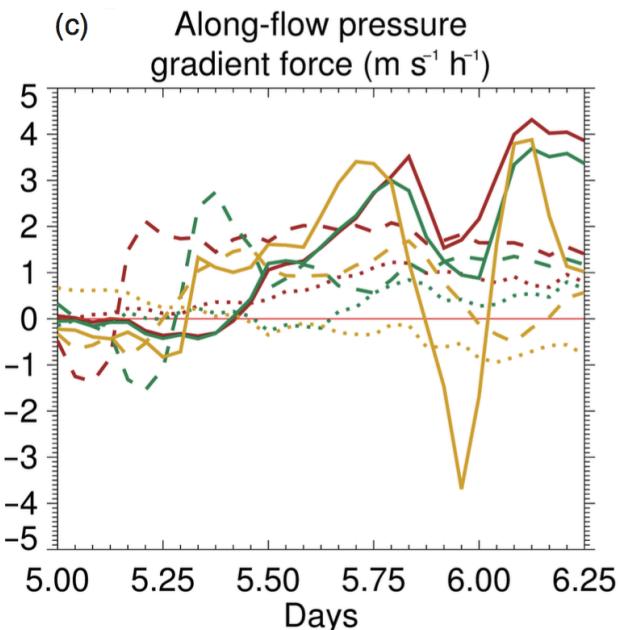
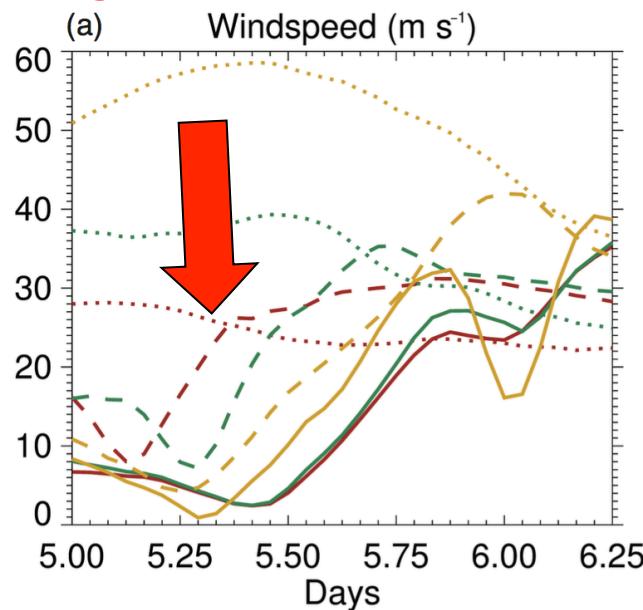


Wind speed and trajectory height from idealized dry baroclinic wave



(Slater et al. 2015, QJRMS)

Wind speed increases as the horizontal pressure gradient force increases.



Cyclonic - A12, B12, D12
 Descending - A22, B22, D22
 Transitional - A20, B20, D20

- D22
- D20
- D12
- B22
- B20
- B12
- A22
- A20
- A12

Frontogenesis (Petterssen 1936)

$$F = \frac{d}{dt} |\nabla_H \theta|,$$

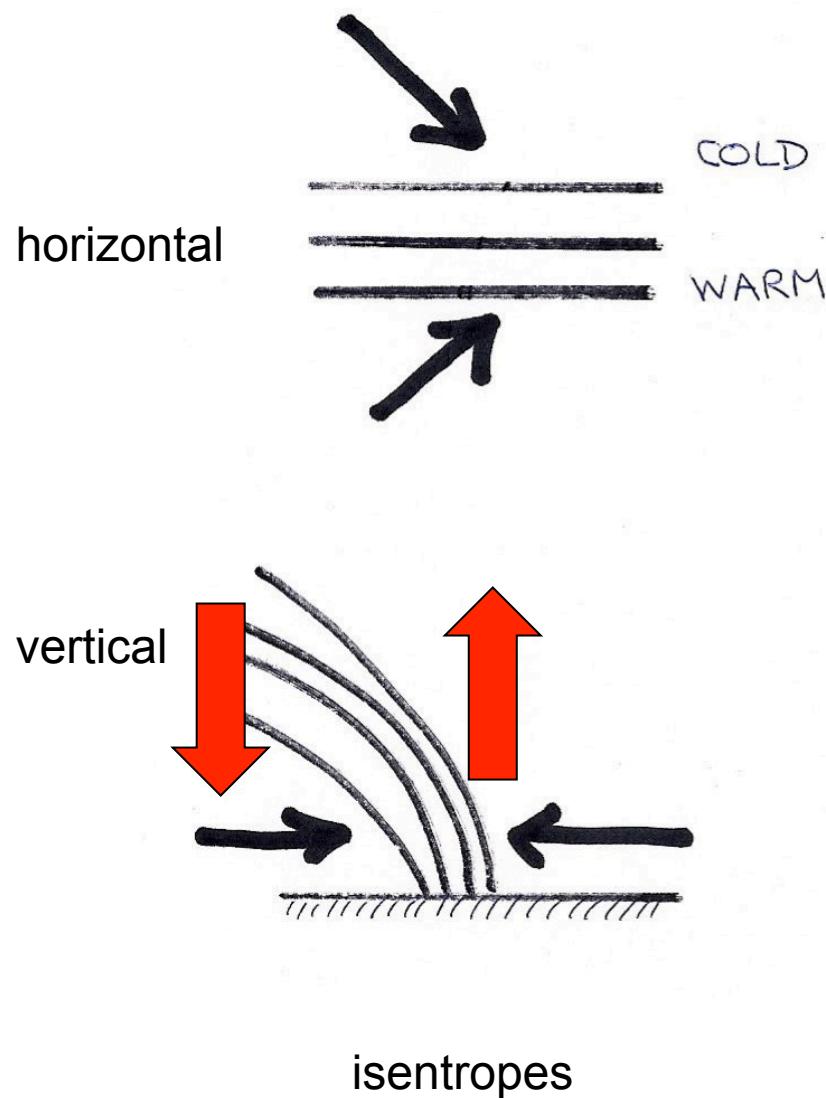
$$\frac{d}{dt} = \frac{\partial}{\partial t} + u \frac{\partial}{\partial x} + v \frac{\partial}{\partial y},$$

$$\mathbf{V}_H = u\mathbf{i} + v\mathbf{j},$$

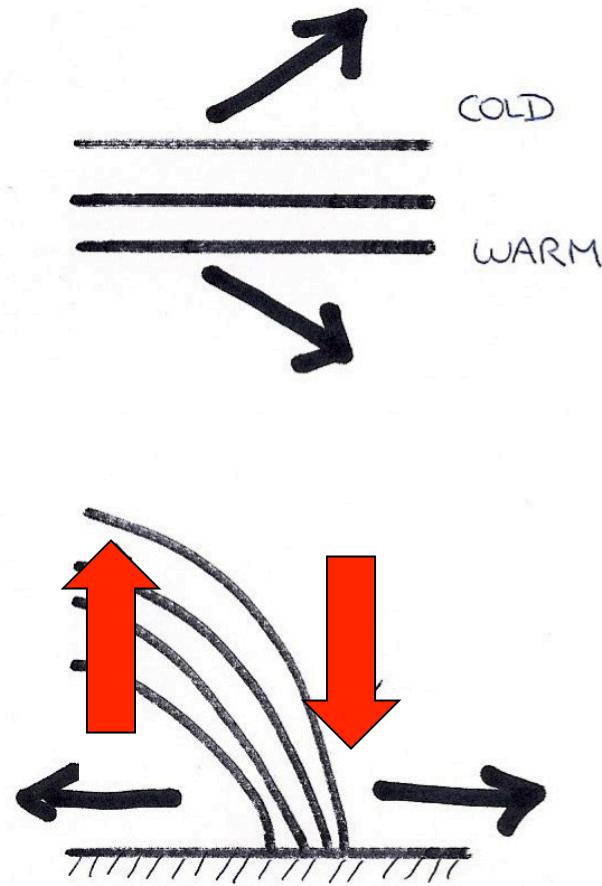
$$\nabla_H = \mathbf{i} \frac{\partial}{\partial x} + \mathbf{j} \frac{\partial}{\partial y}.$$

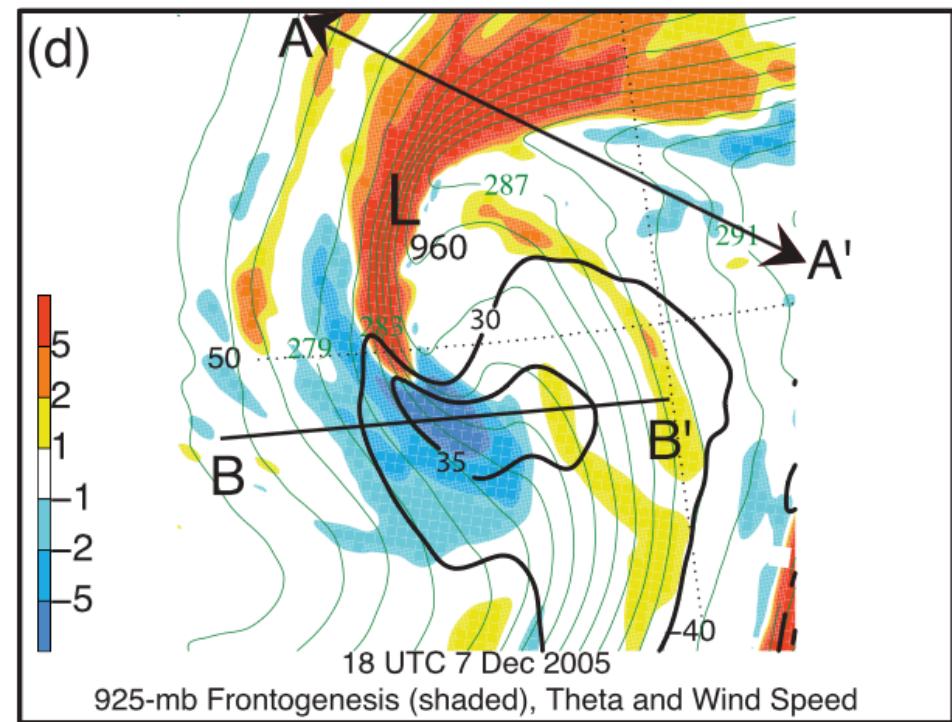
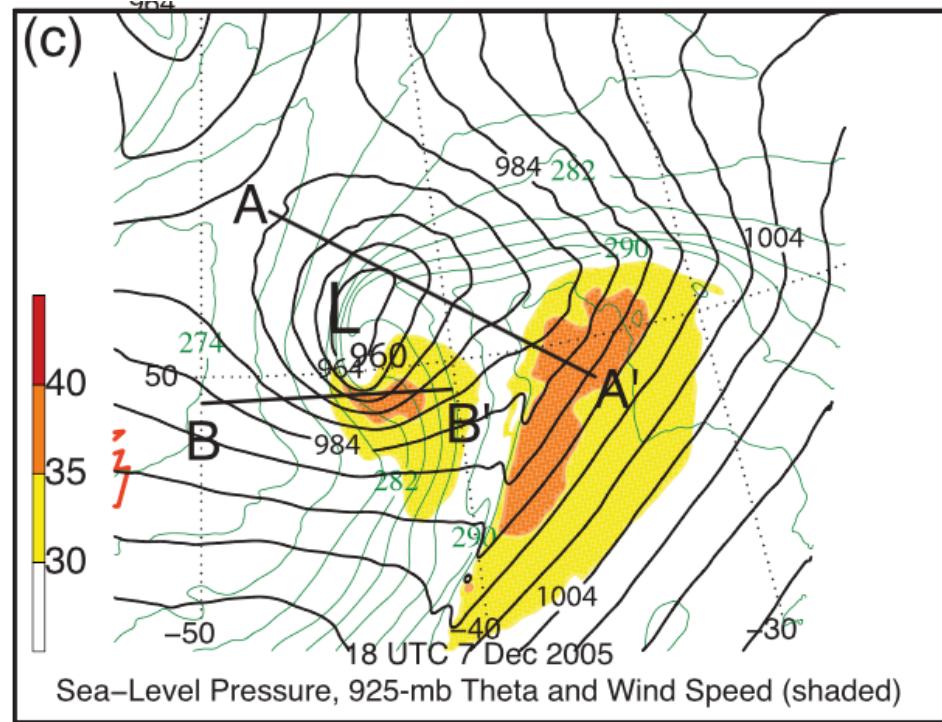
$$F = \frac{1}{2} |\nabla_H \theta| (E \cos 2\beta - \underbrace{\nabla_H \cdot \mathbf{V}_H}_{\text{deformation}} - \underbrace{\nabla_H \cdot \nabla_H}_{\text{divergence}}),$$

Frontogenesis

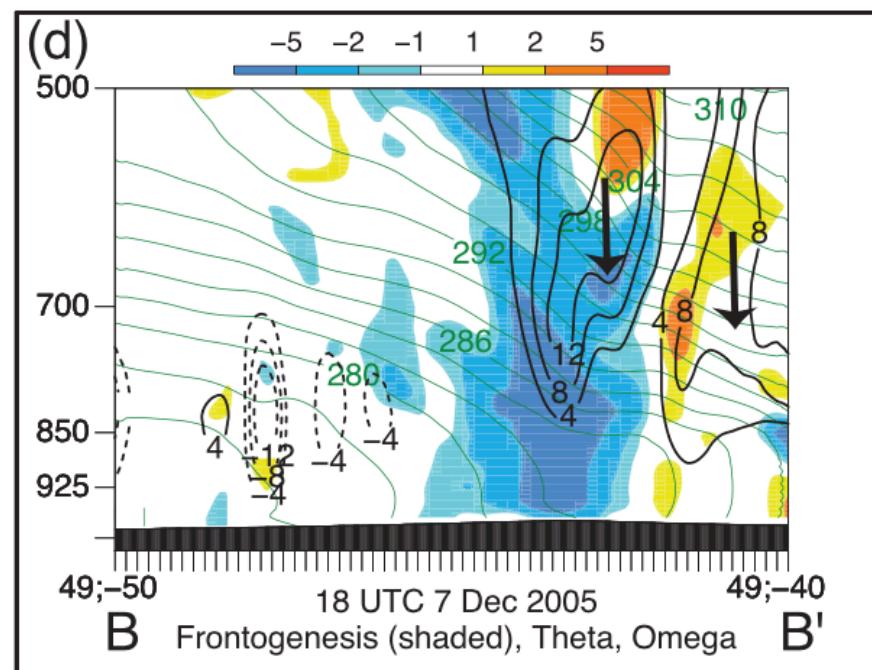
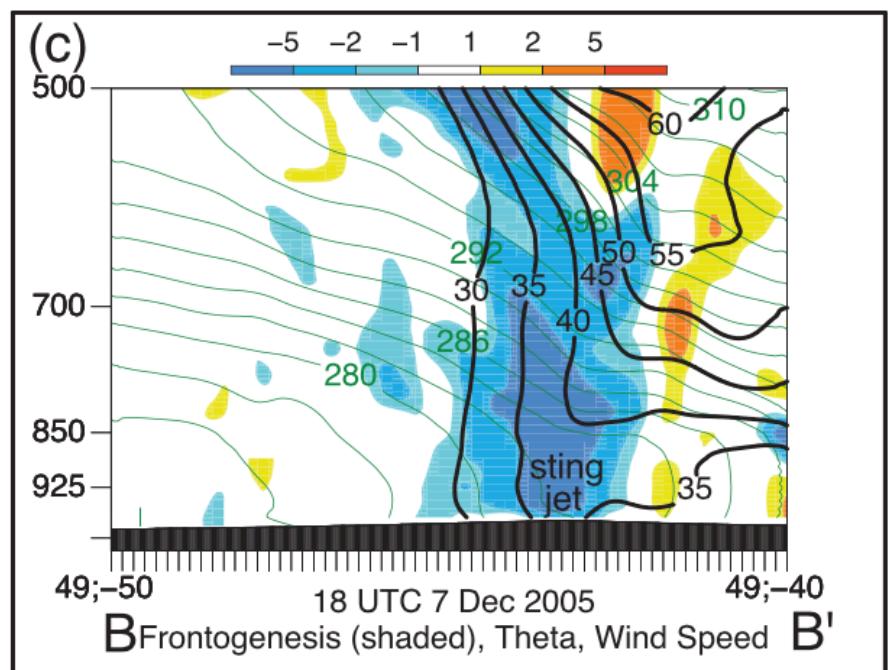
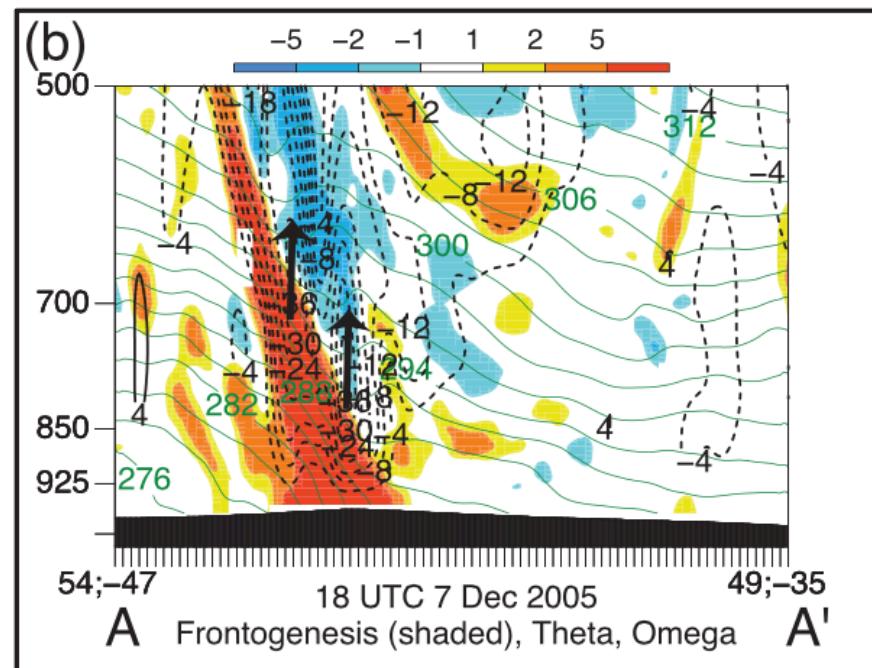
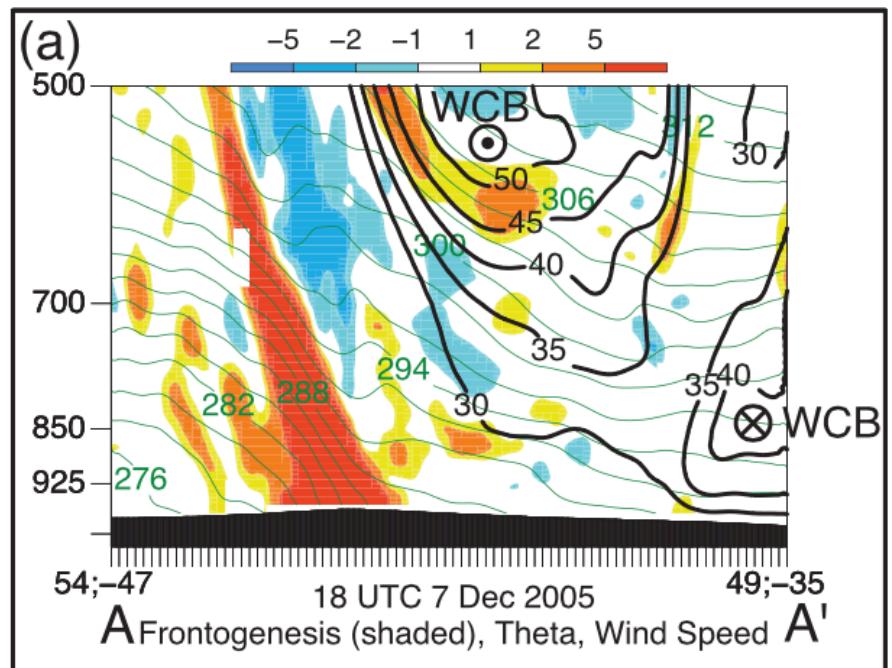


Frontolysis

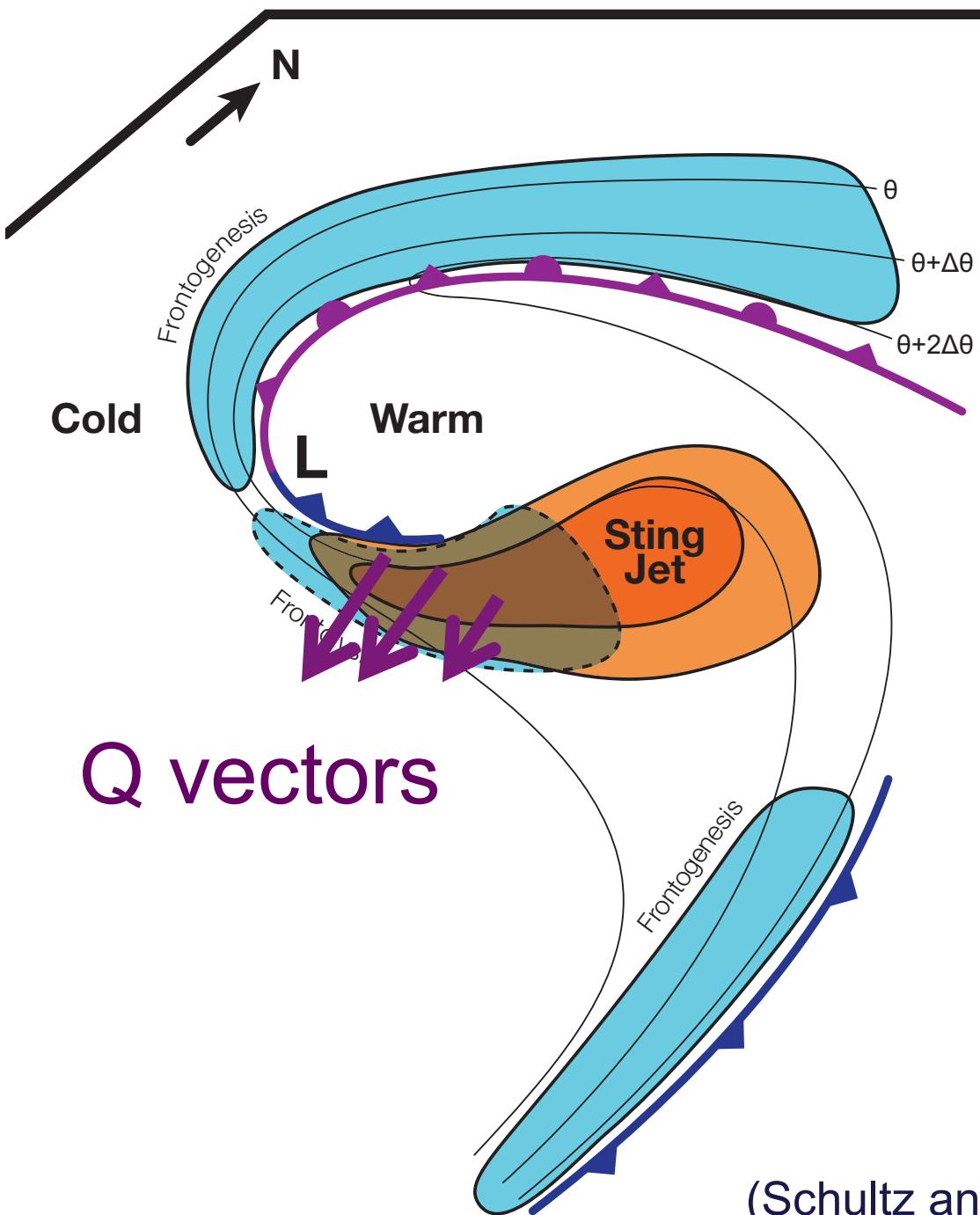




(Schultz and Sienkiewicz 2013)



(Schultz and Sienkiewicz 2013)

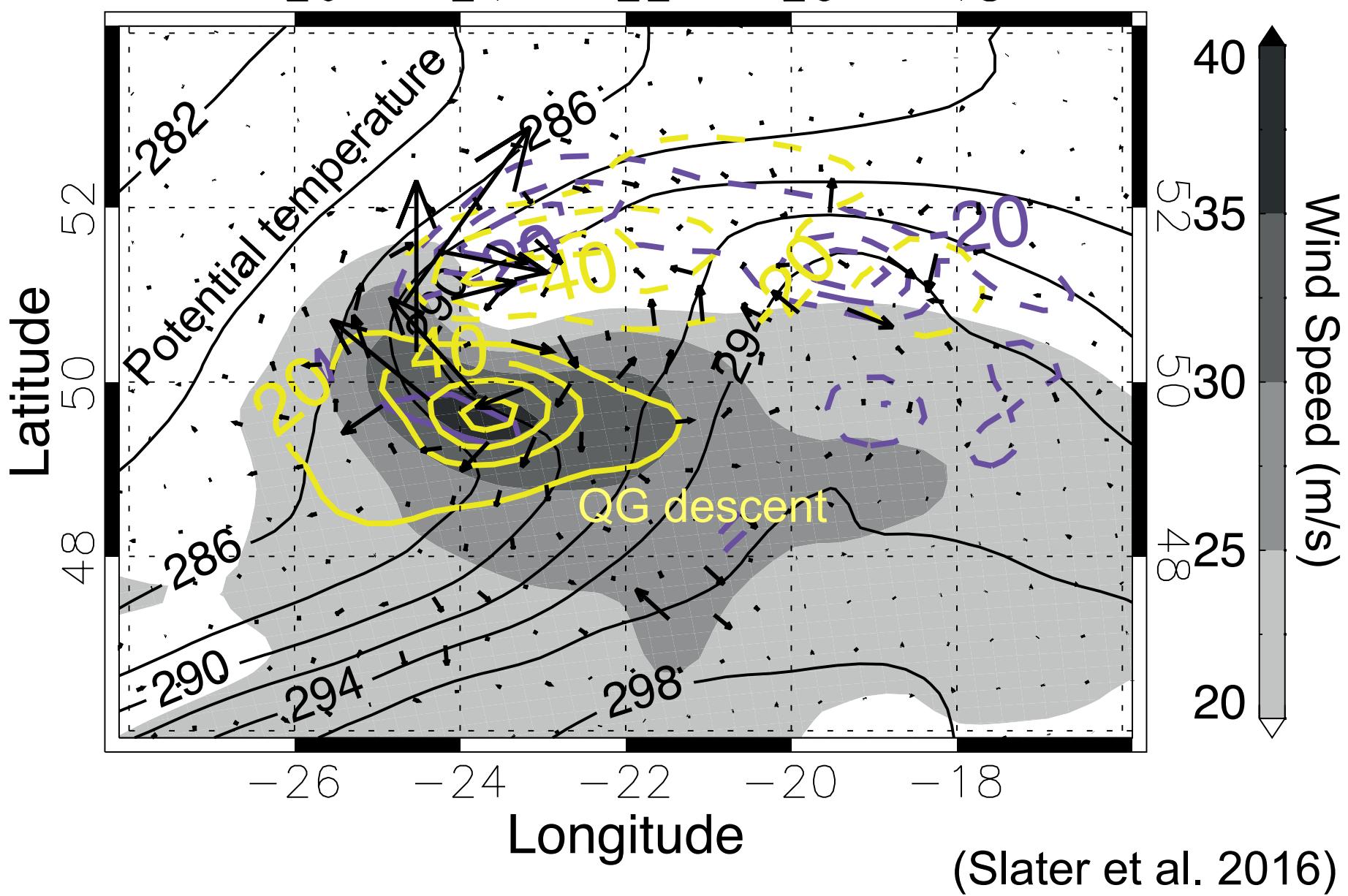


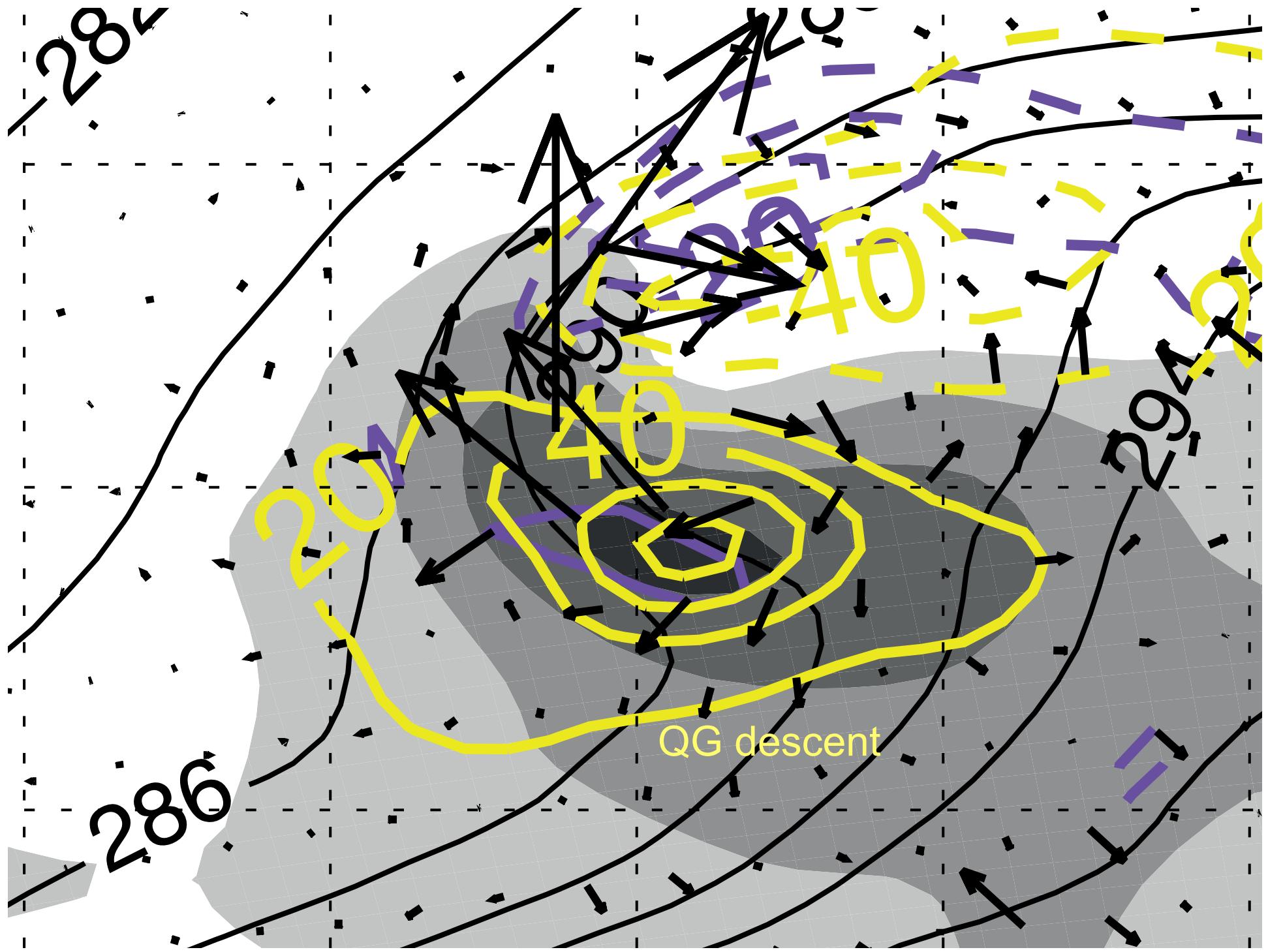
Q vectors

(Schultz and Sienkiewicz 2013)

a) 0100 UTC 12 February CNTL

Q vectors
 $1 \text{ m s}^{-3} \text{ Pa}^{-1}$ →





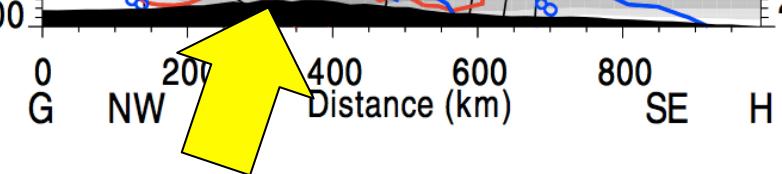
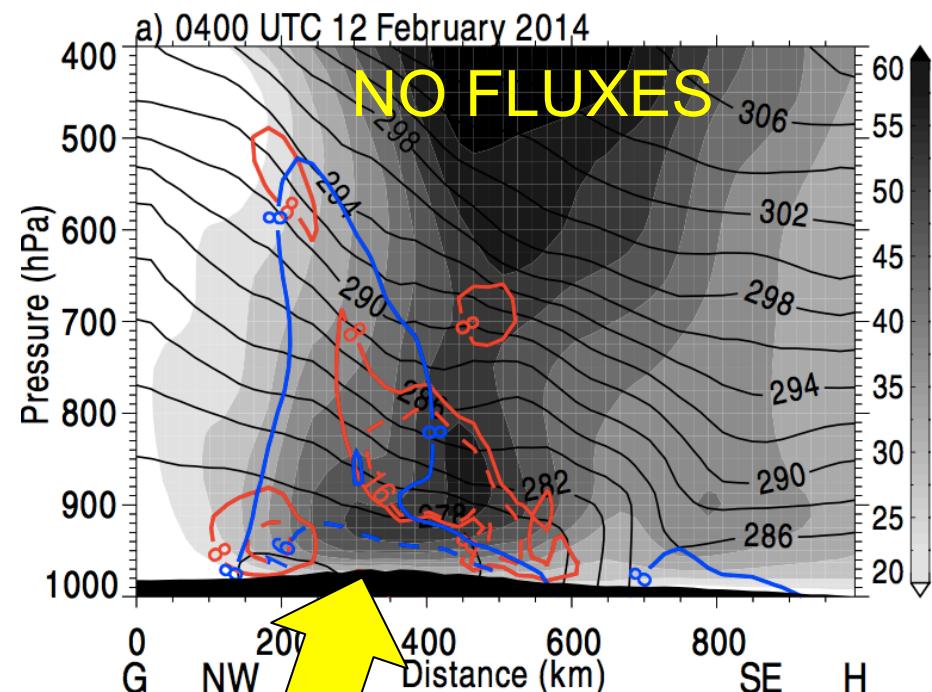
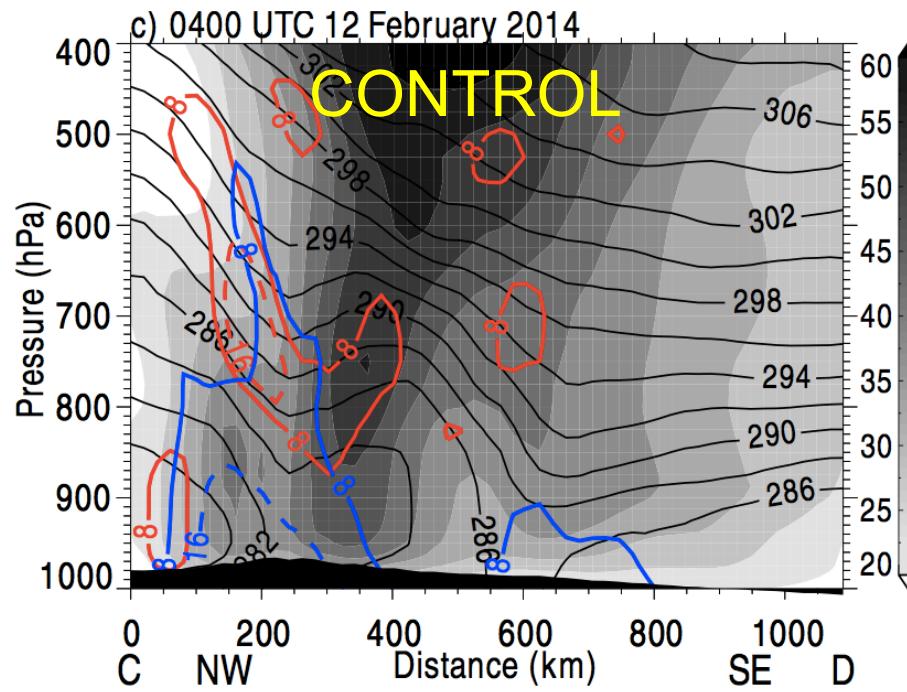
Frontogenesis/frontolysis explains sting jets on the mesoscale.

Why sting jets occur at the end of bent-back front.

Why sting jets occur in Shapiro–Keyser cyclones, but not Norwegian cyclones.

Why trajectories in sting jet descend.

Fluxes from ocean may be critical.



where sting jet
should be

Simulations with no fluxes did not produce strong surface winds.

Slater, Schultz and Vaughan, 2015:
Near-surface strong winds in a marine extratropical cyclone: Acceleration of the winds and the importance of surface fluxes. QJRMS, in revision.