DISCRETE GUST MODEL FOR LAUNCH VEHICLE ASSESSMENTS

Frank B. Leahy*
NASA Marshall Space Flight Center, Huntsville, Alabama

1. INTRODUCTION

Analysis of spacecraft vehicle responses to atmospheric wind gusts during flight is important in the establishment of vehicle design structural requirements and operational capability. Typically, wind gust models can be either a spectral type determined by a random process having a wide range of wavelengths, or a discrete type having a single gust of predetermined magnitude and shape. Classical discrete models used by NASA during the Apollo and Space Shuttle Programs included a 9 m/sec quasi-square-wave gust with variable wavelength from 60 to 300 m (NASA, 2000). A later study derived discrete gust from a military standard (MIL-STD) document that used a “1-cosine” shape (Adelfang and Smith, 1998). The MIL-STD document contains a curve of non-dimensional gust magnitude as a function of non-dimensional gust half-wavelength based on the Dryden spectral model, but fails to list the equation necessary to reproduce the curve (DoD, 1990). Therefore, previous studies could only estimate a value of gust magnitude from the curve, or attempt to fit a function to it (Adelfang and Smith, 1998). Furthermore, the MIL-STD curve is based on a 1% risk gust magnitude, so there was no way to determine gust magnitudes for other risk levels.

This paper presents the development of the MIL-STD curve, and provides the necessary information to calculate discrete gust magnitudes as a function of both gust half-wavelength and the desired probability level of exceeding a specified gust magnitude. The development background is an extension of the descriptions provided in Chalk et al, 1969, and Jones, 1967.

2. BACKGROUND

For analyses of spacecraft vehicles, wind gusts can be treated as either random (spectral turbulence) or discrete. For random gusts, typical spectral models include the Von Karman and Dryden turbulence models. The Von Karman model has widely been considered the more “realistic” model when it comes to defining turbulence spectra. However, due to the computational complexity of the Von Karman model, the Dryden model is typically used in aerospace vehicle analyses. The longitudinal, lateral, and vertical Dryden spectra are:

\[ \Phi_u(\Omega) = \sigma_u^2 \frac{2}{\pi} \frac{\Omega L_u}{1 + (\Omega L_u)^2} \]
\[ \Phi_v(\Omega) = \sigma_v^2 \frac{2}{\pi} \frac{L_v}{1 + (\Omega L_v)^2} \]
\[ \Phi_w(\Omega) = \sigma_w^2 \frac{2}{\pi} \frac{L_w}{1 + (\Omega L_w)^2} \]

where \( \Omega \) is spatial frequency (wavenumber), \( \sigma \) is the turbulence standard deviation, \( L \) is the turbulence scale length, and the subscripts \( u, v, \) and \( w \) denote the longitudinal, lateral, and vertical components, respectively (NASA, 2000). The longitudinal component of turbulence is parallel to the steady-state wind vector, while the lateral and vertical components are perpendicular to it. The non-dimensional Dryden spectra are shown in Figure 1.

The spectra shown herein are in the spatial domain. The spectra can easily be transformed to the frequency domain by use of the Jacobian transform \( \Omega = \omega / V \) where \( \omega \) is radial frequency and \( V \) is the magnitude of the mean wind vector relative to the speed of the aerospace vehicle (NASA, 2000).

An important detail that will be needed later is that the autocorrelation (lag correlation) is determined by taking the inverse Fourier transform of the Dryden spectra. The autocorrelation describes the correlation between gusts separated by a distance or time interval.

Figure 1. Non-dimensional Dryden spectra for the longitudinal, lateral, and vertical components of turbulence.

* Corresponding author address: Frank B. Leahy, NASA/MSFC, Mail Code EV44, Huntsville, AL 35812; e-mail: frank.b.leahy@nasa.gov
For the spatial domain, these autocorrelations are (Hogge, 2004):

\[
R_u(d) = e^{-\frac{d}{2L_u}} \\
R_v(d) = e^{-\frac{d}{2L_v}} \left(1 - \frac{d}{2L_v}\right) \\
R_w(d) = e^{-\frac{d}{2L_w}} \left(1 - \frac{d}{2L_w}\right)
\]

where \( d \) is the lag distance. The autocorrelations as a function of non-dimensional values \( d/L \) are given in Figure 2.

\[d/L\]
\[V_m\]
\[V_m\]
\[d_m\]
\[2d_m\]
\[\text{Distance}\]

Figure 3. Graphical depiction of the 1-cosine discrete gust.

3. DEVELOPMENT OF DISCRETE GUST MAGNITUDES

Generally, random gusts about a mean wind are considered to be normally distributed. This is true for each of the gust components (longitudinal, lateral, and vertical). Therefore, the random gusts, \( V \), will have a probability density function (pdf) of the form:

\[
f(V) = \frac{1}{\sqrt{2\pi\sigma_V^2}} \exp\left(-\frac{1}{2} \left(V - \mu_V\right)^2 / \sigma_V^2\right)
\]

where \( \mu_V \) and \( \sigma_V \) are the mean and standard deviation of the gusts. An initial gust, \( V_1 \), will be related to a gust some distance away, \( V_2 \), by the conditional probability:

\[
f(V_2 \mid V_1) = \frac{f(V_1, V_2)}{f(V_1)}
\]

where \( f(V_1, V_2) \) is the joint pdf of the two gusts. The mean values of the gust, \( V_1 \) and \( V_2 \), are simply the value of the mean wind. The important factor in this development is the deviation of gusts around a mean, not the value of the mean itself. Therefore, these can be removed by setting them equal to zero. If the initial gust, \( V_1 \), is assumed to start at zero, then Equation 5 will reduce to:

\[
f(V_2 \mid V_1 = 0) = \frac{1}{\sqrt{2\pi\sigma_V^2(1 - \rho^2)}} \exp\left(-\frac{1}{2(1 - \rho^2)} \frac{V_2^2}{\sigma_V^2}\right)
\]

where \( \rho \) is the correlation between \( V_1 \) and \( V_2 \). If we let

\[
e = \sigma_V \sqrt{1 - \rho^2}
\]

Figure 2. Autocorrelation as a function of non-dimensional values \( d/L \).

The discrete gust provides a “spike-type” input whose magnitude is based on information from the spectral model. Along with the gust magnitude, the gust gradient over a specified temporal or spatial interval is also important. The classical shape of the discrete gust is that of a “1-cosine” shape given in Equation 3 and shown in Figure 3 as a function of gust width (Chalk et al, 1969).

\[
V = \begin{cases} 
V_m & \text{for } 0 < x < d_m \\
\frac{V_m}{2} \left[1 - \cos\left(\frac{\pi x}{d_m}\right)\right] & \text{for } 0 \leq x \leq 2d_m \\
0 & \text{for } x > 2d_m
\end{cases}
\]

Here, \( V \) is the gust magnitude at distance \( x \), and \( V_m \) is the gust magnitude at \( d_m \), the gust half-width. Several values of \( d_m \) can be chosen to tune the gust width to excite desired vehicle responses. The next section describes the development of the methodology to determine appropriate discrete gust magnitudes for given gust half-widths.
and substitute into Equation 6, we get the common form of the normal pdf (see Equation 4):

$$f(V_2 | V_1 = 0) = \frac{1}{\sqrt{2\pi \sigma^2}} \exp \left( -\frac{1}{2} \frac{V_2^2}{\sigma^2} \right)$$ (8)

The cumulative distribution function (cdf) then becomes:

$$F(V_2) = P(x \leq V_2) = \frac{1}{\sqrt{2\pi \sigma^2}} \int_{-\infty}^{V_2} \exp \left( -\frac{1}{2} \frac{V^2}{\sigma^2} \right) dV$$ (9)

Since Equation 9 can not be integrated in closed form, computer routines are relied upon to determine the normal cdf. To determine the gust magnitude value for $V_2$, computer routines to calculate the inverse of the normal cdf are needed. Here, one would input the probability level ($P$), the mean of the gusts ($\mu = 0$), and the standard deviation ($\sigma$). To determine $\sigma$, substitute the appropriate value of $R$ from Equation 2 for $\rho$ in Equation 7.

Care must be taken when performing this calculation. Recall that gust magnitudes are normally distributed and are two-sided. That is to say, the probability of having a gust magnitude greater than $V_2$ is the same as the probability of having a gust less than $-V_2$. Therefore, the probability level chosen must be the two-sided probability. Most computer routines compute the one-sided probability. To account for this, let $P = 1 - (risk/2)$, where risk is the desired probability of exceeding a given gust magnitude.

Appropriate turbulence standard deviations, $\sigma$, and length scales, $L$, are provided in Table 2-79b of NASA, 2000 for various altitudes and turbulent cases (light, moderate, and severe). Table 1 below lists the values from 1 to 10 kilometers for the severe case.

Table 1. Severe turbulence standard deviations and length scales for the longitudinal (U), lateral (V), and vertical (W) turbulent components.

<table>
<thead>
<tr>
<th>Altitude (km)</th>
<th>Standard Deviation (m/sec)</th>
<th>Length Scale (m)</th>
</tr>
</thead>
<tbody>
<tr>
<td></td>
<td>U</td>
<td>V</td>
</tr>
<tr>
<td>1</td>
<td>5.70</td>
<td>4.67</td>
</tr>
<tr>
<td>2</td>
<td>5.80</td>
<td>4.75</td>
</tr>
<tr>
<td>4</td>
<td>6.24</td>
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<td>8</td>
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<td>5.98</td>
</tr>
<tr>
<td>10</td>
<td>7.72</td>
<td>6.00</td>
</tr>
</tbody>
</table>

4. EXAMPLE CALCULATION AND COMPARISON TO PREVIOUS STUDIES

As an example, a longitudinal gust magnitude for severe turbulence at the 10 km level for a 500 m gust half-width is computed. The turbulence standard deviation and length scale at this level are $\sigma = 7.72$ m/s and $L = 1230$ m. The risk of exceeding the gust magnitude is 1%. Therefore, the probability to input into the inverse normal computer routine is $P = 1 - (0.01/2) = 0.995$. The resulting gust magnitude is 14.83 m/sec. Figure 4 shows the longitudinal gust magnitude for this example with gust half-widths in the range 1 to 1000 m.

![Figure 4. Longitudinal gust magnitude for varying gust half-width, $\sigma = 7.72$ m/sec and $L = 1230$ m.](image)

A previous study (Adelfang and Smith, 1998) attempted to fit a function that reproduced the MIL-STD curve, which is non-dimensional gust half-width ($d_m/L$) vs. non-dimensional gust magnitude ($V_m/\sigma$). The two curves are shown in Figure 5. Notice that the Adelfang and Smith curve is slightly conservative compared to the MIL-STD curve. Figure 6 is the same as Figure 4, but with the corresponding Adelfang and Smith curve included. The percent difference between the two curves is depicted in Figure 7. The Adelfang and Smith curve over estimate the MIL-STD curve by as much as 13% in the 200 to 300 m gust half-width range.

5. SUMMARY

A development method of the MIL-STD non-dimensional discrete gust magnitude curve has been presented. The development allows for slight reduction in gust magnitudes determined by other studies (Adelfang and Smith, 1998). Also, the new technique allows for selection of various risk levels of exceeding gust magnitudes (see Figure 8).

6. ACKNOWLEDGEMENTS

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7. REFERENCES


