# P6R.5 SYSTEMATIC VARIATIONS OF RAINDROP SIZE SPECTRA WITH ALTITUDE DERIVED FROM WIND PROFILER: MEASUREMENTS FOR TRMM PR EVALUATION

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# INTRODUCTION

Measurements of raindrop size distribution (DSD) are crucial for understanding rain formation processes in which the rainfall rate changes significantly, corresponding to the size spectrum (Atlas et al., 1999, 2000). Many processes affect DSD as drops fall, i.e., coalescence, nucleation, diffusion growth or evaporation, sedimentation, and breakup. Numerical modeling results show significant changes in the DSD during rainfall (e.g., List and McFarquhar 1990). However, few observations of the vertical profile of DSD are reported (Kobayashi and Adachi, 2001). More observations are needed.

This paper presents vertical variations of raindrop size distributions in precipitation occurred associated with typhoon 9707, on 20 June 1997 in Tukuba, Japan. A 404 MHz wind profiler (Fig.1) is used to derive DSD of precipitation at the Meteorological Research Institute, Tsukuba, Japan where is located about 100 km east of the typhoon center. The rain rate was ranging from 1 to 30 mm/h at the ground. It was continuously observed by the wind profiler. Table 1 summarizes the wind profiler system parameters

# DSD RETRIEVAL FROM WIND PROFILER

UHF/VHF wind profilers detect signals scattered from turbulent fluctuations in the atmospheric refractive index, known as Bragg scattering (Williams et al. 1995). Received signals, which are Doppler shifted by air motion, are used to measure wind profiles. Under some conditions, the wind profiler can also detect signals from hydrometeors by Rayleigh scattering. A vertically pointing wind profiler can also measure the fall velocity of raindrops. Raindrop size can be derived from the fall velocity using a relationship between raindrop size and the terminal velocity. However, measured fall velocity combines the terminal velocity of raindrops and the ambient air motion. Small errors in the estimates of vertical air speed and turbulence can lead to significant errors in the derived DSD. Air motion should therefore be removed from the measured fall velocity.

\**Corresponding author address:* T. Kobayashi, Meteorological Research Institute, Tsukuba, 305-0052, Japan Email: kobay@mri-jma.go.jp Doppler spectra measured with a wind profiler often have two distinct spectra, one from air motion and one from precipitation. Under such conditions, the wind profiler can measure both the fall velocity of raindrops and the ambient air motion, including the mean speed and the turbulence. The two distinct spectra facilitate the removal from the spectra of air motion effects.

In this paper, we have derived DSD by using a method developed by Kobayashi and Adachi (2005). The method is based on an iterative procedure and can derive arbitrarily shaped DSD, named ITRAN (ITerative Retrieval method for Arbitrarily shaped N (D): raindrop size distribution). This method assumes no prior shape of drop size distributions and fully derives automatically; additionally, it can be applied to large data volumes. Furthermore, it is insensitive to initial values. Here, we will briefly explain the retrieval method.



Fig.1 A photograph of the MRI wind profiler.

# Table 1: MRI Wind profiler system parametersFrequency: 404.0 MHzPower: 35 kW (peak), 1.5 kW (average)Pulse width: 1.67, 6.67 $\mu$ secondsPRT: 100 or 153.5 $\mu$ secondsFirst range gate 500 m (minimum)Antenna size: 10.4X10.4 m (phased array)Antenna gain: 33 dBiBeamwidth: $\leq$ 4.1°Beam steerability: 0 - 360° (azimuth), 0 - 15°(zenith)

The spectral density ( $S_{obs}$ ) measured with a wind profiler is given by ,

$$S_{obs}(v) = [P_D S_D(v - v_0) * S_t(v) + P_t S_t(v - v_0)] * W,$$
(1)

where *v* is the Doppler velocity (upward +) and  $v_0$  the mean vertical air velocity.  $P_t$  and  $P_D$  are the echo powers associated with refractive index irregularities and precipitation, respectively.  $S_t$  and  $S_D$  are the normalized spectral densities due to air turbulence and precipitation, respectively. *W* is the window function. The asterisk represents the convolution operation. The precipitation spectrum  $S_D$  is related to the raindrop size distribution N(D) (*D*: raindrop diameter) as

$$S_D(v) = D^0 N(D) dD / dv / Z,$$
(2)

where Z is the reflectivity factor. Here we consider only spectral broadening due to turbulent motion. When N(D) is derived, broadening effects in a measured Doppler spectrum should be removed by deconvolution.

The turbulent spectrum is assumed to have a Gaussian distribution, given by

$$S_t(v) = \frac{1}{(2\pi\sigma^2)^{1/2}} \exp(-v^2/2\sigma^2), \qquad (4)$$

where  $\sigma$  is the spectral width. The Gaussian parameters are determined by non-linear least squares fitting using a Marquart method.

The turbulent peak velocity ( $v_{max}$ ) is easily identified if the measured spectrum has two distinct peaks. The boundary of the two peaks ( $v_{min}$ ) is the minimum position or is a few FFT points apart from the  $S_t$  peak. The spectrum ranging from  $v_{max}$ - $v_{min}$  to  $v_{max}$ + $v_{min}$  is fitted. If  $S_t$  is correctly positioned, the fitting works well.

Next, an iterative deconvolution of  $S_D^*S_t$  is applied.

$$S_d^{i+1}(v) = \langle S_d^i(v) \rangle \cdot \left\langle \frac{S(v)}{S^i(v)} \right\rangle_{i=0\dots\infty},$$
(5)

where  $X^i$  is the value of  $X (S_d \text{ or } S)$  at the <sup>1</sup>th iteration.  $S_d^i$  is defined as  $P_D S_D$  in (1).  $S^i(v)$  is the sum of  $P_t S_t$ (v) and the convolution of  $S_d^i * S_t$ . The symbol < > denotes a running mean. The initial estimate  $S_d^0$  is the observed spectrum smoothed by 1 ~ 5-points running mean, depending on the degree of statistical fluctuation. Multiplying  $S_d^i$  by the ratio of the observed spectrum S (v) to  $S^i(v)$  yields a new estimate of  $S_d^{i+1}$ . This process, by which new estimates of  $S_d^{i+1}$  are substitute in, is repeated. The precipitation spectrum is obtained point by point. Thus any shape of the DSD can be derived.

### OBSERVATIONS

Corrected Doppler spectra measured with a 404 MHz profiler are available for 5 hours from 11:00 to 16: 00 JST at time intervals of 6 minutes and altitude intervals of 250 m. Doppler spectra of a low mode vertically pointing beam were averaged for one minute. Doppler spectra of a distinct peak were selected.

Here we will focus on a relationship between the rain water content (RWC) and the median volume diameter (D<sub>0</sub>). Numerical models for the raindrop formation involving breakup and coalescence processes suggest that there should be correlations between the rain rate and the typical drop size. Many studies have reported the statistical relations between various integral parameters like the radar reflectivity factor, total number of drops (N<sub>T</sub>), etc. derived from surface disdrometer measurements. However, no clear correlations have been reported between the rain rate and the median volume diameter.

Figure 2 shows wind profiler-derived  $D_0$  versus the rain rate (R). The integral parameters were calculated from the derived DSD for diameters ranging from 0.5 to 5.5 mm to avoid from the effects of different drop size range because the minimum raindrop diameter of the derived DSD varied from 0.1 to 0.5, depending on atmospheric and rain conditions. Figure 2 shows no apparent correlation between R and  $D_0$ , as in studies by Bringi et al. (2002). Testud et al. (2001) described aircraft measurements in the Tropical Ocean and Global Atmosphere Coupled Ocean-Atmosphere Response Experiment and showed weak correlation for the entire dataset. He also showed an inverse relationship between N\* (intercept value of DSD) and D<sub>0</sub>. During fall of raindrops, small and large drops, respectively, increase and decrease in number when collision-



Fig. 2: Scatter plot of the wind profiler-derived median volume diameter versus the rain rate.

induced breakup occurs. These changes cause  $D_0$  to decrease and  $N_T$  to increase. The opposite also occurs through the collision-induced coalescence. Therefore, an inverse relation between  $N_T$  and  $D_0$  is expected. However, not even a weak relation between R to  $D_0$  is present in Fig. 2 despite the apparent relation between R and  $N_T$  (not shown here). Median volume diameter characterizes the shape of DSD but is mathematically independent on  $N_T$ . Similar values of  $D_0$  are observed during precipitation events with similar DSD and different R. This may be the reason of no clear correlation between R to  $D_0$  shown in Fig.2.

Figure 3 shows a scatter plot of  $D_0$  versus rain water content normalized by  $N_{\rm T}$ . A clear relation is present, which is in agreement with results from Testud et al. (2001). The relationship between  $N_{\rm T}$  and  $D_0$  may be smeared when many events are considered. For a single precipitation event,  $N_{\rm T}$  can be expected to relate to  $D_0$ .

Figures 4, 5, and 6 show vertical profiles of RWC,  $D_0$  and  $N_T$  at 11:50, 11:57, and 13:09. The rain rates were 1, 6, and 10 mm/h, respectively at the ground. The values of RWC and  $N_{\rm T}$  were normalized by the maximum values in the profile. The RWC profiles have peaks and dips at altitudes corresponding to those of the  $N_{\rm T}$ . On the other hand, contrast to the uncorrelated results shown in Fig.2, clear inverse relations between  $D_0$  and  $N_T$  appear. The  $N_T$  have peaks at 3.3 km at which the  $D_0$  is minimum value in Fig.4. In Fig.5, a significant increase in the RWC at altitude of 2 km corresponds to a rapid decrease in  $D_0$ . Although most changes in RWC are in phase with change in  $N_{\rm T}$ , there are some examples of inverse relationships between  $N_T$  and RWC. For example, at 4 km, RWC increases but  $N_{\rm T}$  decreases in both figures 4 and 6. At this altitude,  $D_0$  increases significantly, which leads to an increase in RWC. Fitting errors in



Fig.3 Scatter plot of D<sub>0</sub> and RWC/NT.









Fig.6 Same as Fig.4 except at 13:09.

the clear-air vertical motion overestimate  $D_0$  and underestimate RWC and vice versa. However,  $D_0$  is fairly insensitive to errors in  $v_0$ . The vertical air speed was almost zero. The vertical profile of Z is similar to that of RWC as estimated from the linear correlation. Errors in  $D_0$  and RWC ( $N_T$ ) are about 10 and 25percent even for an error of 0.5 m/s in  $v_0$ . The effects can be negligible associated with much larger variations of  $D_0$  and  $N_T$ .

# CONCLUTIONS

Precipitation associated with typhoon 9707 was observed by using a 404 MHz wind profiler on 20 June 1997 in Tsukuba, Japan. Vertical variations of the rain drop size distributions were derived by using an iterative retrieval method. This technique assumes no prior shape of DSD and therefore enables the vertical variations of  $D_0$  as well as rain rate to be analyzed in detail. The median volume diameter is statistically independent of the RWC. The large variability of the total number of drops  $N_T$  between the rain events leads to the uncorrelated results. A vertical profile of each rain event shows a clear inverse relationship between  $N_T$  to  $D_0$ .

Whether these interesting results reported here are seen in other area or not is a problem to be solved. The precipitation radar (PR) onboard the Tropical Rainfall Measuring Mission (TRMM) has measured the vertical profile of rain rate since 1997. The rainfall rate is estimated from the radar reflectivity factor by using a power law:  $R=aZ^b$  in which the parameter *a* and *b* are closely related to both N<sub>T</sub> and D<sub>0</sub>. Even for the current system, in which a single-frequency radar is operating, we can estimate raindrop size to some degree. For the Global Precipitation Mission (GPM), in which dual-frequency radar will be equipped, we can estimate the raindrop size more accurately, and therefore, we will examine the interesting results reported here for entire global scale.

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